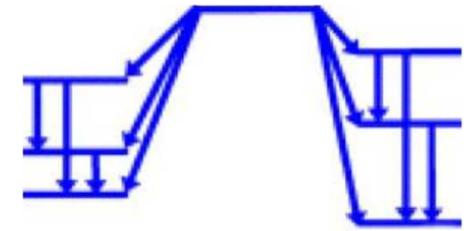


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ICRM GSWG

GAMMA-SPECTROMETRY

INTRODUCTION

General characteristics

Energy resolution

Peak area

Attenuation

Corrective factors

Gamma-ray spectrometry is a fast and widespread method

- used to identify radionuclides and
- to determine their activity.

- Non-destructive method
- Simple sample preparation
- But relative method (requires calibration using standards)

Mature technique

- Button-push method (**which software ?**)
- Hidden difficulties / uncertainties factors
- Evolution
 - Large detectors
 - Digital electronics
 - Compton-suppression

RESULTS DEPENDING ON EQUIPMENT CHARACTERISTICS :

HPGe detector (other detectors for low-energy X-rays (Si(Li), SDD), room temperature detectors NaI(Tl), CdTe, CZT...), etc.)

Sample (material, shape)

Measurement conditions

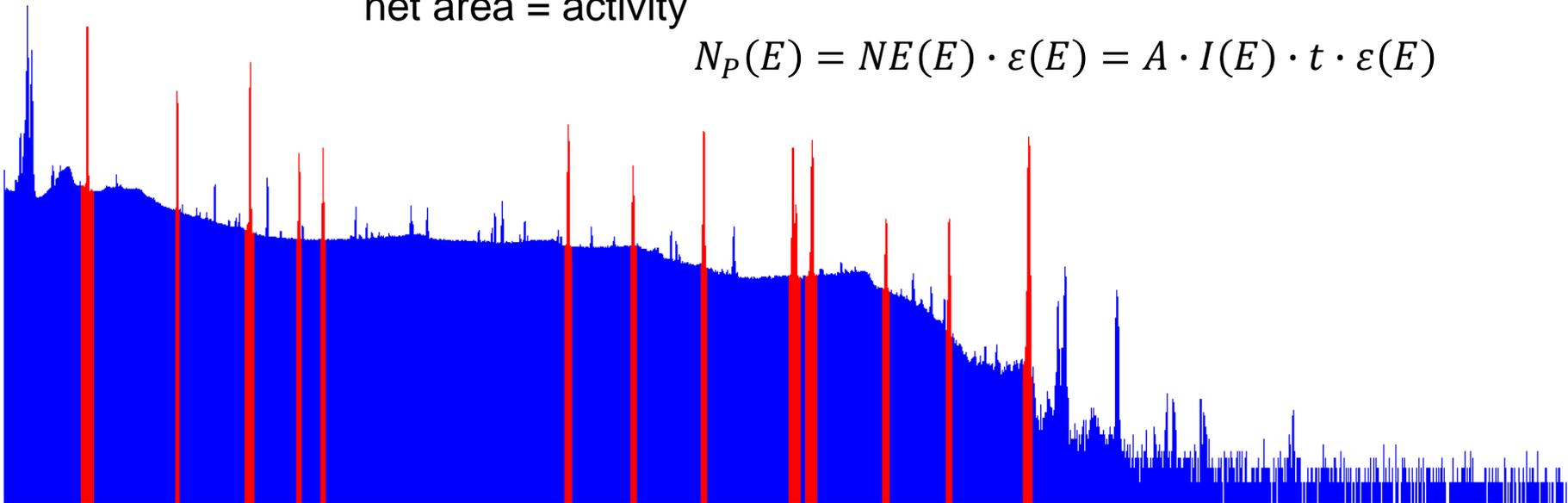
^{152}Eu – Point source at 10 cm from the detector window (100 cm³) HPGe detector

Peaks (full-energy peaks) :

position = energy -> radionuclide identification;

net area = activity

$$N_p(E) = NE(E) \cdot \varepsilon(E) = A \cdot I(E) \cdot t \cdot \varepsilon(E)$$



Other features (continuum) : partial energy deposition, scattering, escape ...

$$N_p(E) = NE(E) \cdot \varepsilon(E) = A \cdot I(E) \cdot t \cdot \varepsilon(E)$$

or simply:

$$n_i = A \cdot I_i \cdot \varepsilon_i$$

A : activity (Bq)

n_i : net peak area

I_i : photon emission intensity

ε_i : detection efficiency

-> 3 results depending on available input parameters:

$$\varepsilon_i = \frac{n_i}{A \cdot I_i}$$

$$A = \frac{n_i}{I_i \cdot \varepsilon_i}$$

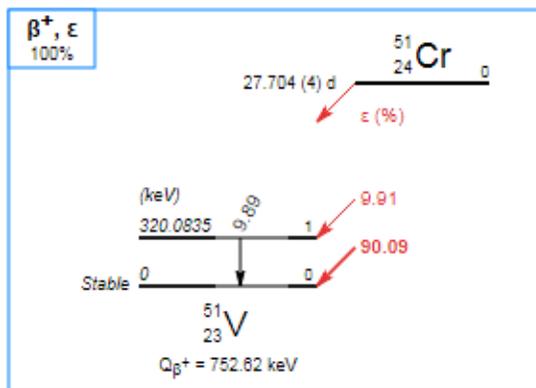
$$I_i = \frac{n_i}{A \cdot \varepsilon_i}$$

Hidden difficulties : peak area, efficiency for any energy (see “Fitting” section), corrective factors

- Decay scheme: Difference emission/transition
- Gamma transition = **gamma emission** + internal conversion

$$\longrightarrow T_g = I_g + I_C$$

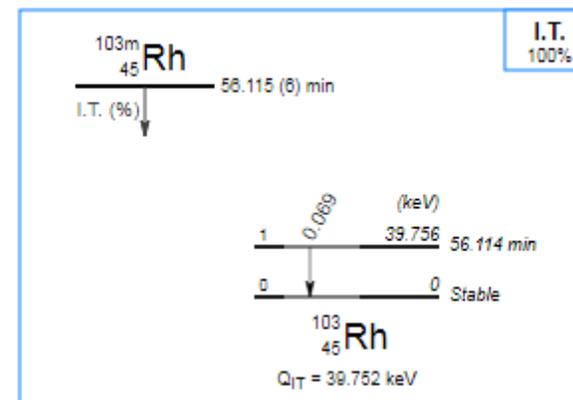
- Conversion coefficient : $\alpha_T = \frac{I_\gamma}{T_\gamma} \longrightarrow I_\gamma = \frac{T_\gamma}{(1 + \alpha_T)}$



$$^{51}\text{Cr} - T_\gamma = 9.91(2) ; I_\gamma = 9.89(2)$$

$$\alpha_T = 0.00181$$

$$I_\gamma = \frac{T_\gamma}{(1 + \alpha_T)} = \frac{9.91}{1.00181} = 9.89$$



$$^{103m}\text{Rh} - T_\gamma = 100 ; I_\gamma = 0.069(4)$$

$$\alpha_T = 1448$$

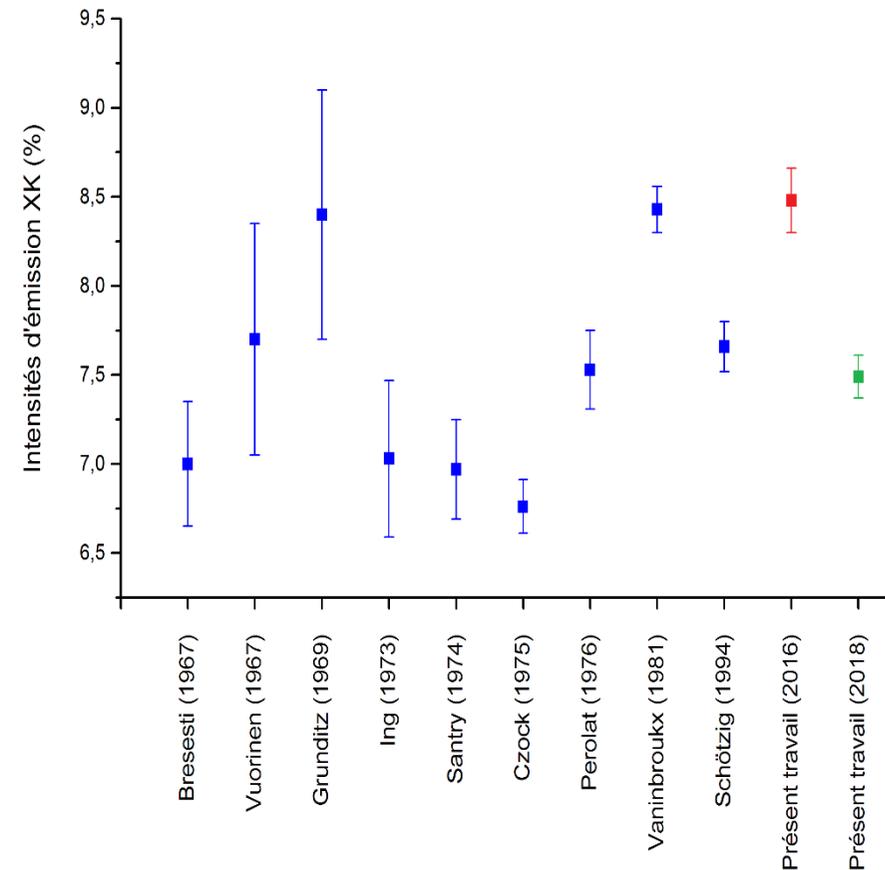
$$I_\gamma = \frac{T_\gamma}{(1 + \alpha_T)} = \frac{100}{1449} = 6.9$$

RADIOACTIVE DECAY DATA

Example:

^{103m}Rh X-ray emission intensities
(short half-life)

Difficulty to derive « recommended » data
requires carefully detailed information

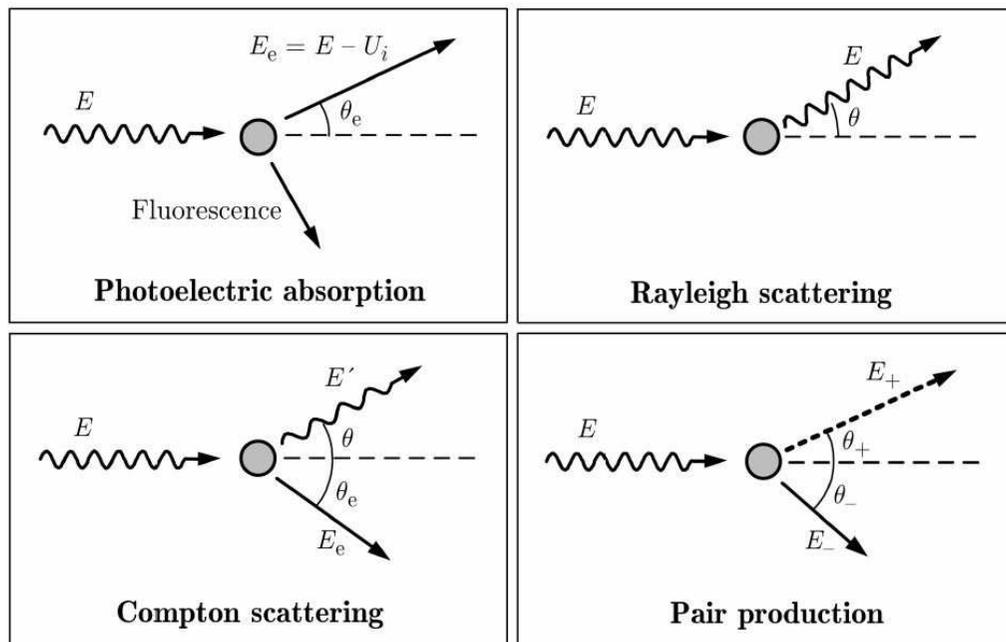


In October 2005, the General Meeting of the *International Committee for Radionuclide Metrology* (ICRM) formally approved the recommendation made by the Nuclear Data Working Group of using the DDEP evaluated decay data in all future nuclear data studies.

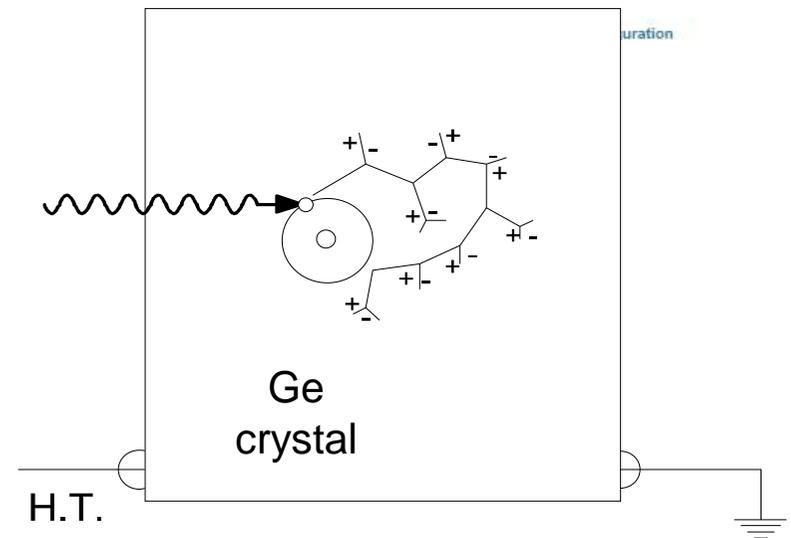
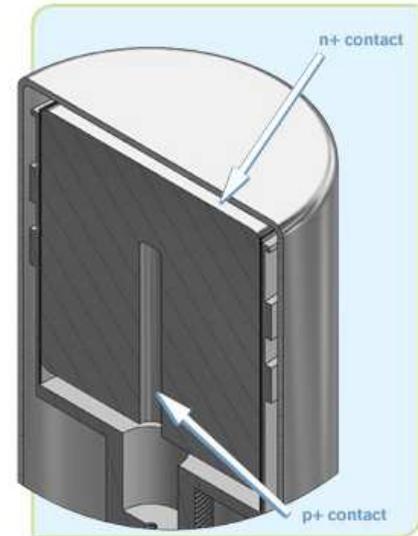
Size -> Efficiency

P or N (low energies)

Photon interaction -> movement of electrons -> electron-hole pairs

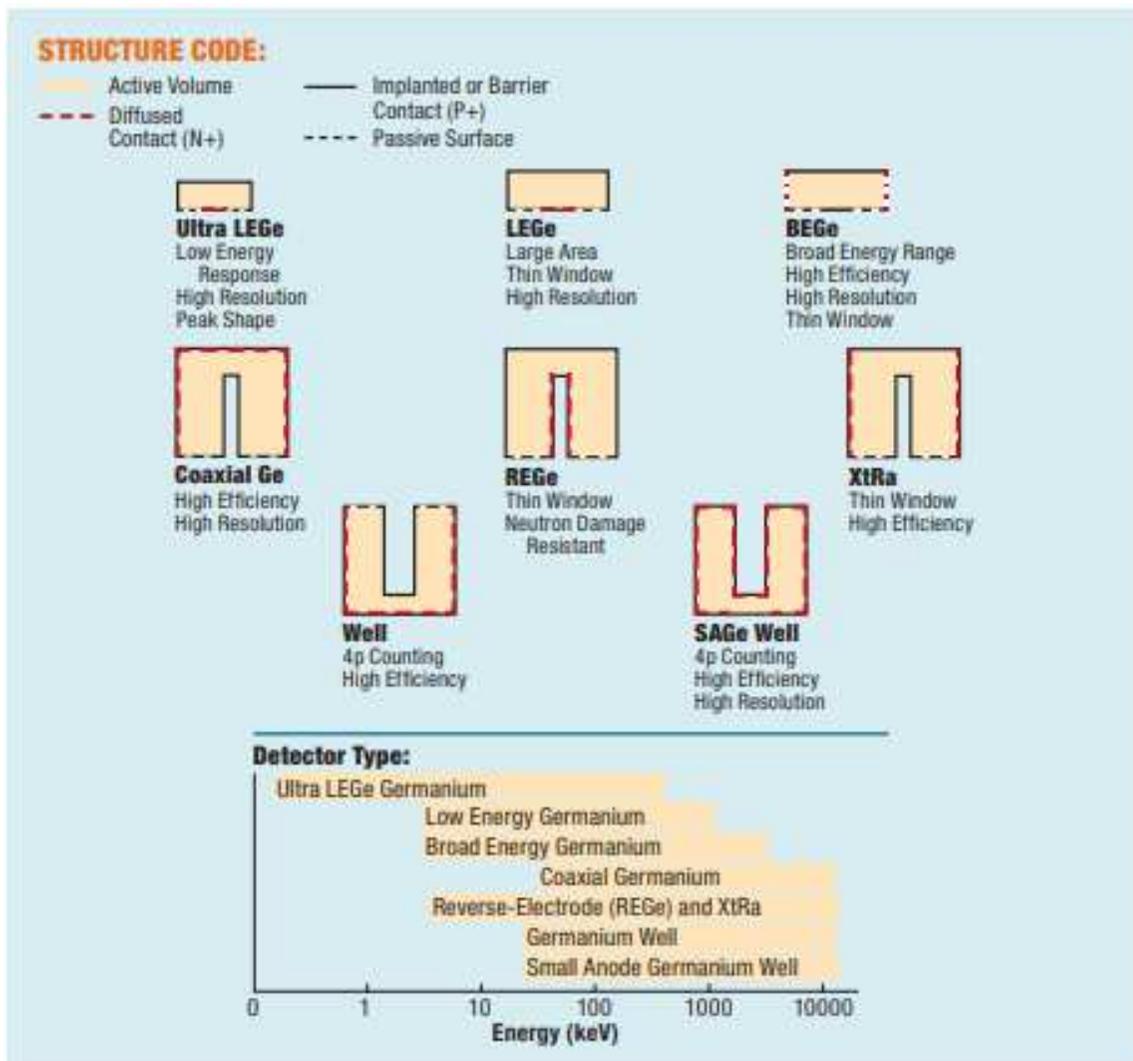


$m_e c^2 \simeq 511 \text{ keV}$, electron rest energy



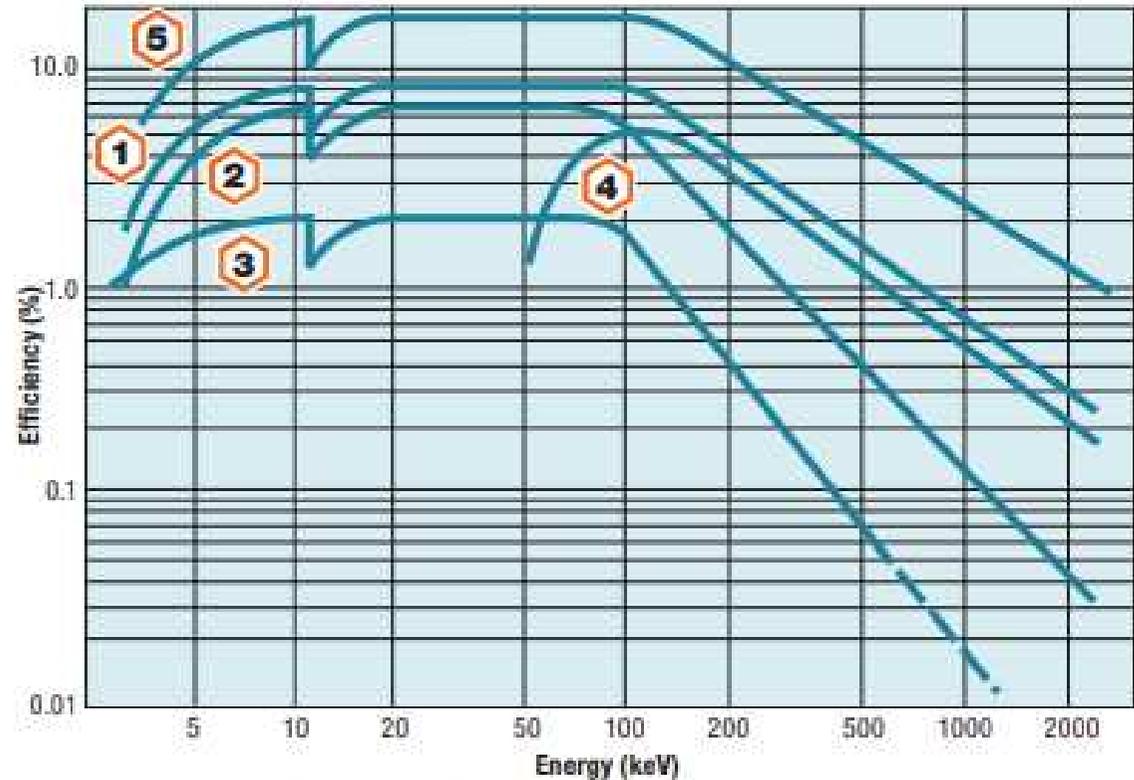
HPGE DETECTOR

Different shapes, types and sizes



HPGE DETECTOR

Different shapes, types and sizes -> influence on detection efficiency



« Relative efficiency »
compared to 3" X 3" NaI (TI)

For 1332 keV (^{60}Co at 25 cm)
(FEP efficiency = $1.2 \cdot 10^{-3}$)

Typical absolute efficiency curves for various Ge detectors with 2.5 cm source to end-cap spacing

- | | |
|---|---|
| 1 REGe, 15% Relative Efficiency Xtra,
15% Relative Efficiency | 3 LEGe, 200 mm ² x 10 mm thick |
| 2 LEGe, 10 cm ² x 15 mm thick | 4 Coaxial Ge, 10% Relative Efficiency |
| | 5 BEGe, 5000 mm ² x 30 mm thick |

Detector + electronic modules (analog/digital)

Preamplifier

- RC
- Reset

Amplifier

- Adapt the shape of the pulse – time constant versus count rate

Multi-channel Analyzer

- Number of channels: 10^3 to $1.6 \cdot 10^4$

It is important to characterize the whole spectrometry chain (detector + electronics)

General characteristics

Energy resolution

Peak area

Attenuation

Corrective factors

Energy resolution :

ΔE (keV) : FWHM : Full Width at Half Maximum of a peak

- Depends on the energy
- Varies from 0.16 keV (@6 keV) to 2.5 keV (@1,3 MeV), depending on the detector

Resolving power : $\Delta E / E$ (%) : Some 10^{-3}
(5 – 10 % for NaI(Tl))

Characteristic of the spectrometer quality :

- separation of closely spaced peaks
- detection limit

ENERGY RESOLUTION

Initial photon emission : **Monoenergetic** line (source) ($\Delta E = \text{some meV}$)

Widening and distortion of the resulting peak (spectrum)

➤ Widening : Gaussian effects:

- ✓ fluctuation of the number of charge carriers : ΔE_s
- ✓ electronic noise : ΔE_E

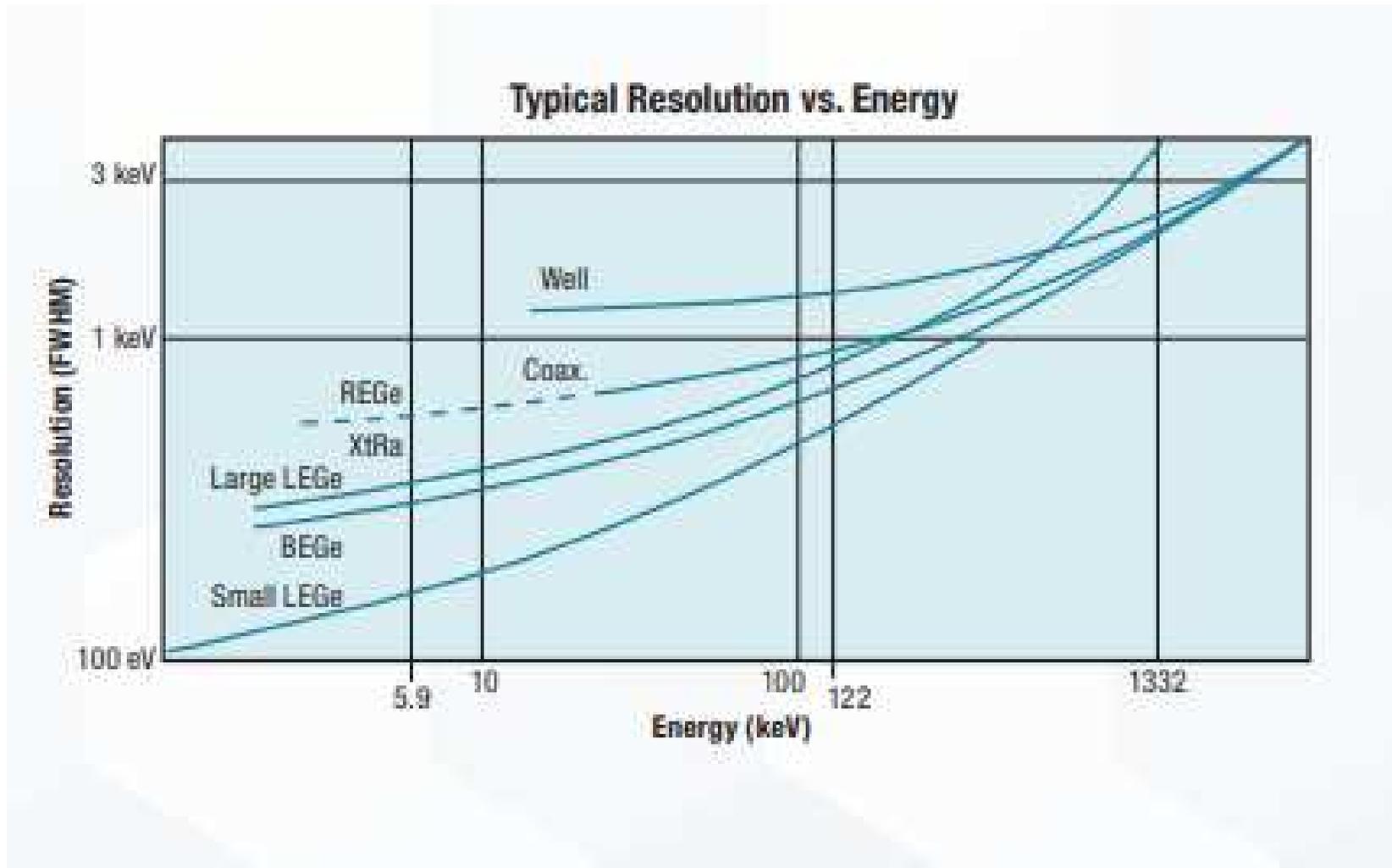
➤ Non Gaussian effects:

- ✓ Charges collection : ΔE_C
- ✓ Drift of the operating system (detector + electronics)
- ✓ Pile-up

Result : in the resulting spectrum : '**peak**' with finite width (**keV**) and more or less symmetrical shape

First approximation : **Gaussian** shape

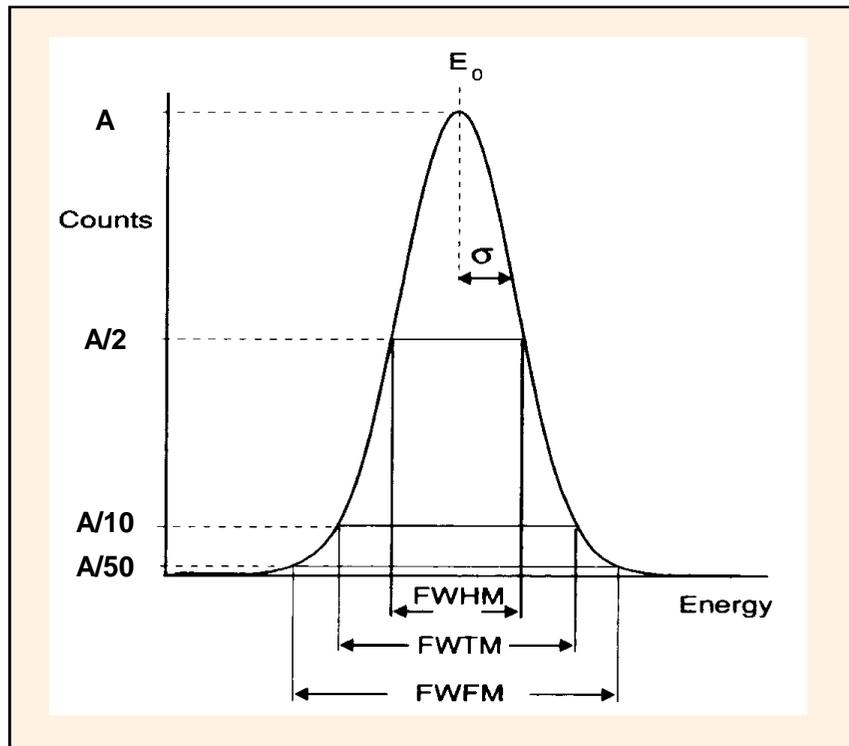
ENERGY RESOLUTION



GAUSSIAN DISTRIBUTION

Distribution centered on energy E_0 with standard deviation, σ

Characteristics :



Full width at half maximum of the Gaussian

$$\text{FWHM } (\Delta E) = 2,355 \sigma$$

$$\text{Area : } S = (2 \pi)^{1/2} \sigma A$$

$$G(E) = A \cdot \exp\left(-\frac{(E - E_0)^2}{2\sigma^2}\right)$$

STATISTICAL NOISE : ΔE_S

Cause : fluctuations of the number of charge carriers n

E : incident energy (some keV)

w : mean pair creation energy (some eV)

$$n = \frac{E}{w}$$

Hypothesis : Poisson statistics : $\sigma_n = \sqrt{n}$

Standard deviation of the number of charge carriers

-> Standard deviation of the deposited energy ($n w$)

$$\sigma_{ES} = w \cdot \sigma_n = w \cdot \sqrt{n} = w \cdot \sqrt{\frac{E}{w}} = \sqrt{w \cdot E}$$

STATISTICAL NOISE : ΔE_S

Example : (at 77K : $w = 2.96$ eV in Ge
 $w = 3.76$ eV in Si)

$E = 1$ MeV \rightarrow in Ge : $n = 3.3 \cdot 10^5$

$$\sigma_{ES} = \sqrt{w \cdot E} = \sqrt{2.96 \cdot 10^6} = 1.72 \text{ keV}$$

$$\Delta_{ES} = 2.355 \cdot \sigma_{ES} = 4.05 \text{ keV}$$

???

The observed width is much lower (< 2 keV)

$\rightarrow F = \text{Fano factor} : F = \frac{\text{Observed variance}}{\text{Poisson predicted variance}}$

STATISTICAL NOISE : ΔE_S

$$\Delta_{ES} = 2.355 \cdot \sqrt{F \cdot \sigma_{ES}^2} = 2.355 \cdot \sqrt{F \cdot w \cdot E}$$

Fano factor < 1 (charges creation are not independent : correlations)

Measured for Ge and Si : value depending on the quality of the crystal ?

Experimental values : (increasing versus time ...)

Ge : 0.07 to 0.12

Si : 0.08 to 0.12

Preamplifier – amplifier

- Independent on the energy
- Can be determined using a pulser : input at the preamplifier « test »

Causes : trapping of charge carriers and ballistic deficit

Trapping (impurities or crystal imperfections) : loss of charge or slowing of the rate of charge collection

Ballistic deficit (electron-hole mobility (15% difference in Ge))

Depends on the quality of the crystal and on the electric field (position of the interaction in the detector)

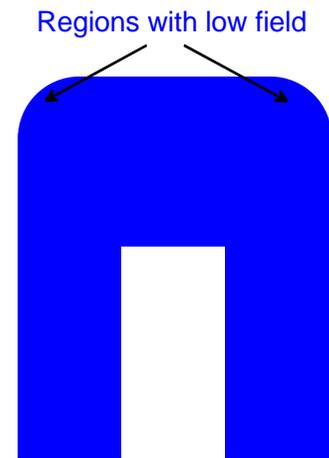
Consequence of the loss of charge :
Low-energy tailing

Improvement :

Increasing the voltage (recommended value by supplier)

Rejecting pulses with slow rise-time

Using collimator to avoid interaction in low electric field regions (front corners)



$$\Delta E^2 = \Delta E_S^2(E) + \Delta E_C^2 + \Delta E_E^2$$

ΔE_S = statistical noise (dependent on the energy)

ΔE_C = charge collection

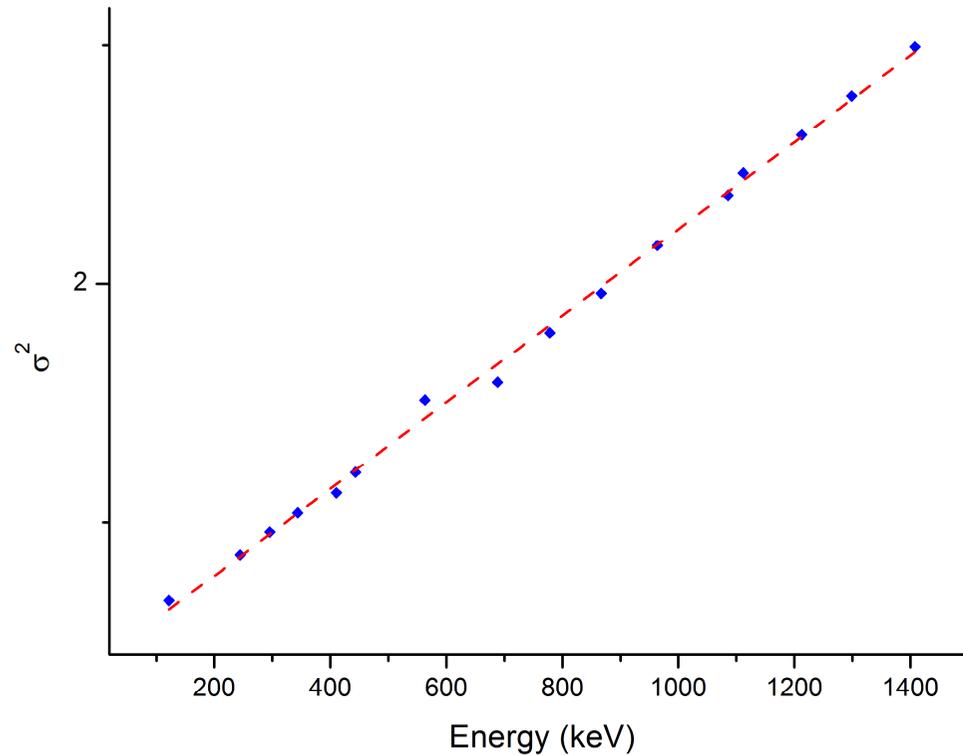
ΔE_E = electronical noise

First approximation : $\Delta E^2 = K + \Delta E_S^2(E)$

$$\Delta E^2 = K + 2.355^2(F \cdot w \cdot E)$$

FWHM versus the energy

$$\Delta E^2 = K + 2.355^2(F \cdot w \cdot E)$$



Linear fitting : $\text{FWHM}^2 = 0.00182E + 0.4141$
 $K = 0.414$ keV (electronics + charge collection)

$$(2.355)^2 F w = 0.00182 \rightarrow F = 0.110$$

FWTM (Full Width at Tenth Maximum)

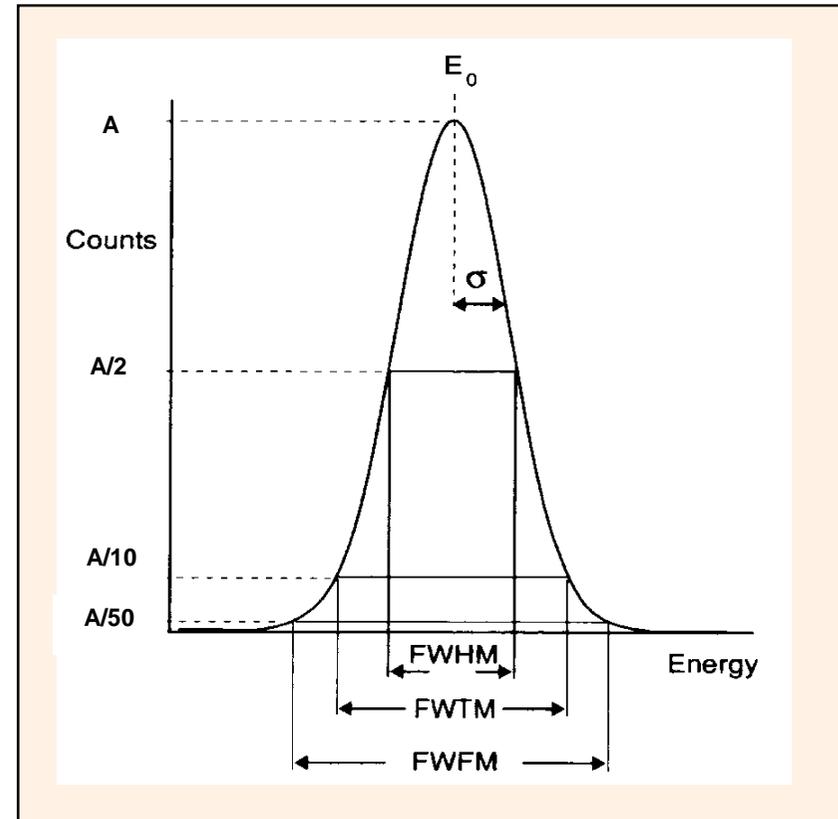
FWFM (Full Width at Fiftieth Maximum)

These values can be compared to those characteristic of a Gaussian (ideal ?) peak :

$$\text{FWTM} / \text{FWHM} = 1.82$$

$$\text{FWFM} / \text{FWHM} = 2.38$$

Information about the peak shape, particularly on its base (collection defects)



General characteristics

Energy resolution

Peak area

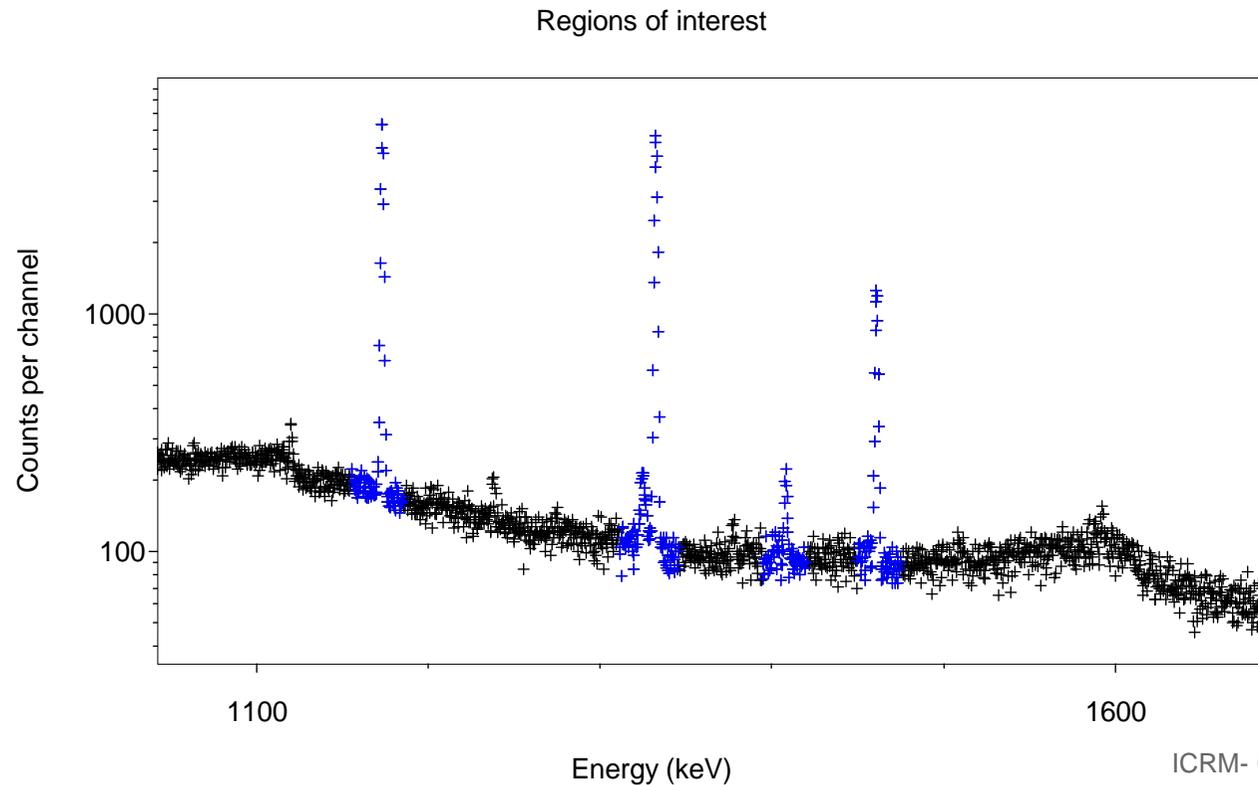
Attenuation

Corrective factors

For quantitative analysis : determination of $N(E) = \text{peak area}$

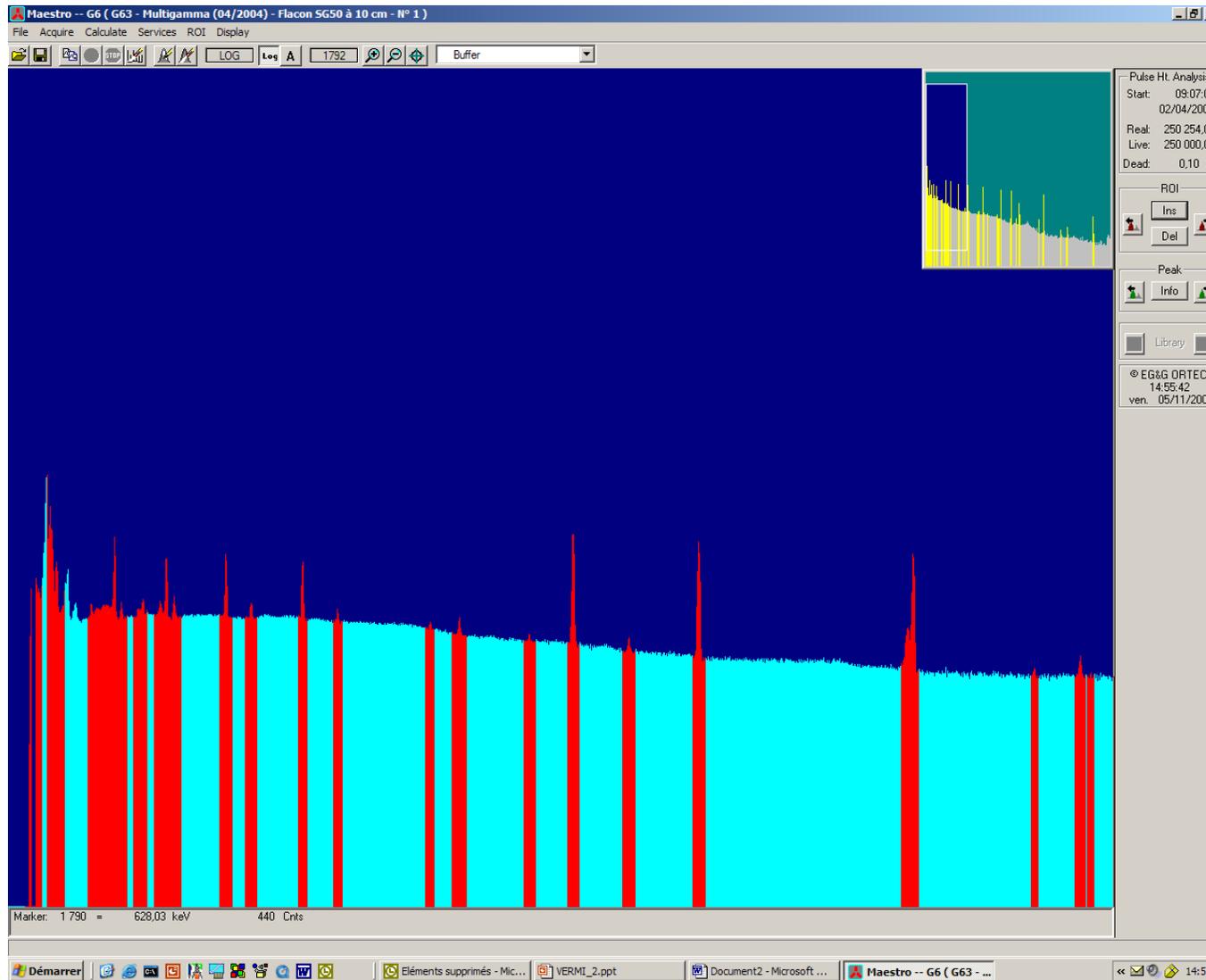
(number of photons with energy E that deposited all their energy in the detector)

Spectrum - > "Regions Of Interest" (ROI's) containing one or several peaks



FINDING PEAKS IN SPECTRUM

Automatic peak search



AUTOMATIC PEAK SEARCH

Peak region :

$$Y(E) = G(E) + (aE + b)$$

(Gaussian + linear background)

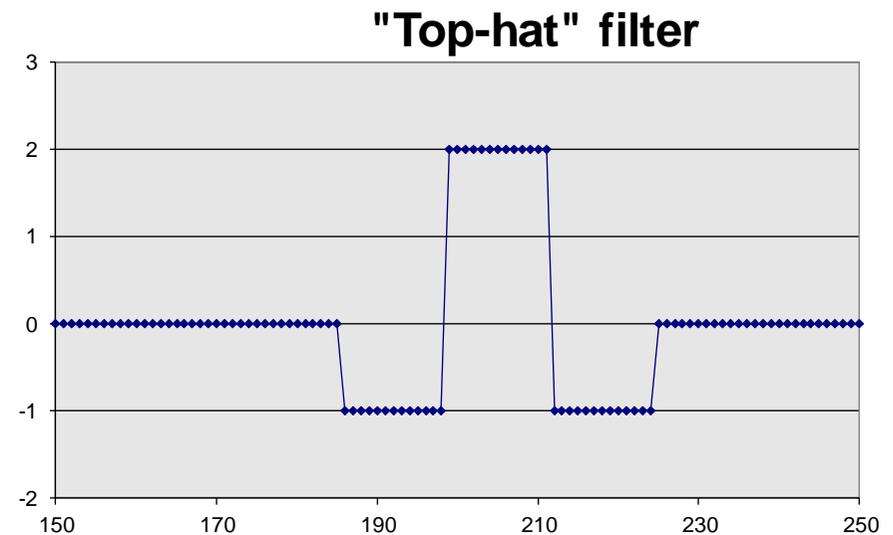
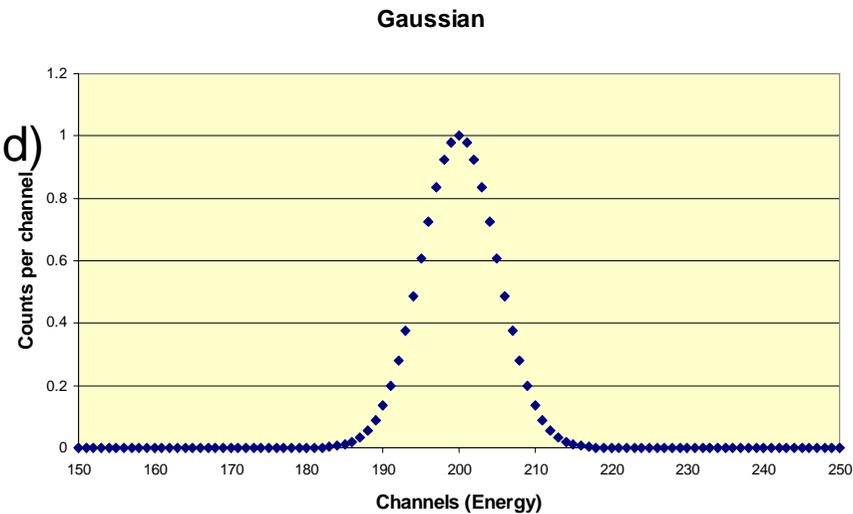
Convolution by a « top-hat » filter

With zero area

Width = Gaussian width

(requires resolution calibration)

$$Y^C = Y \otimes TF$$

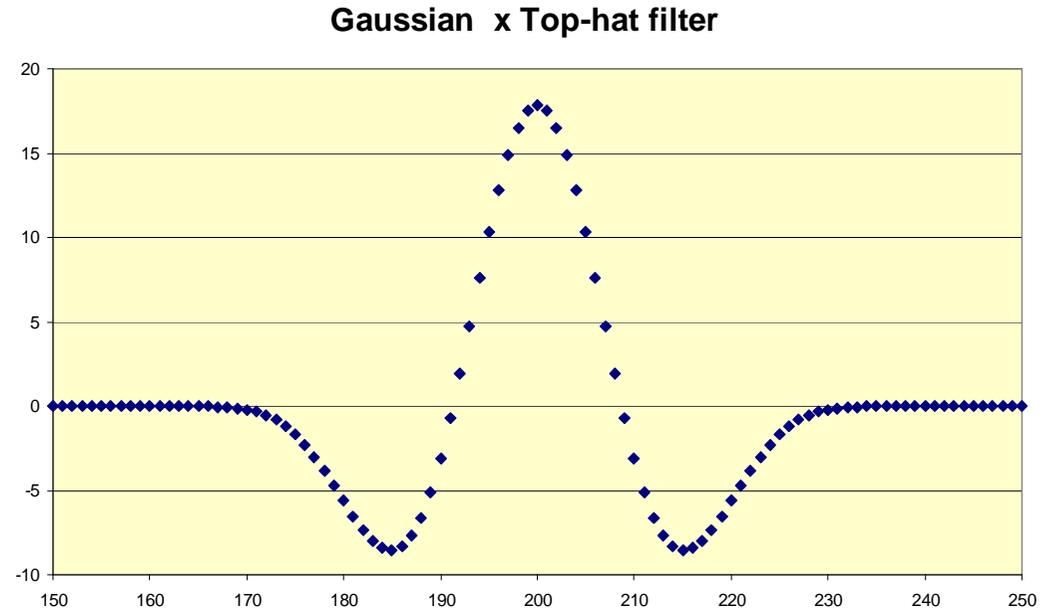


AUTOMATIC PEAK SEARCH

- ◆ Presence of a peak when $Y^C > T$ (threshold)

Peak FWHM \leftrightarrow channels
where $Y^C = 0$

\rightarrow definition of a ROI



Peak detection sensitivity is dependent on the threshold value :

T too low \rightarrow spurious peaks

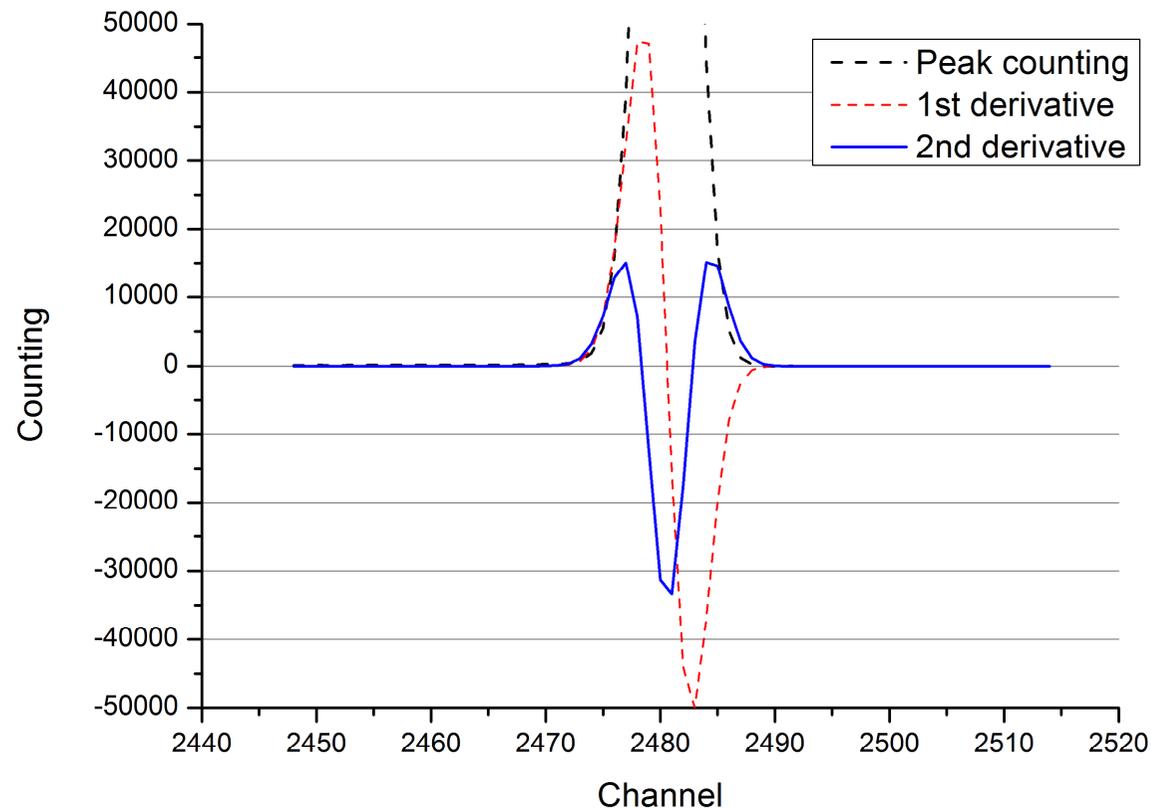
T too high \rightarrow missing peaks

AUTOMATIC PEAK SEARCH

DERIVATIVES

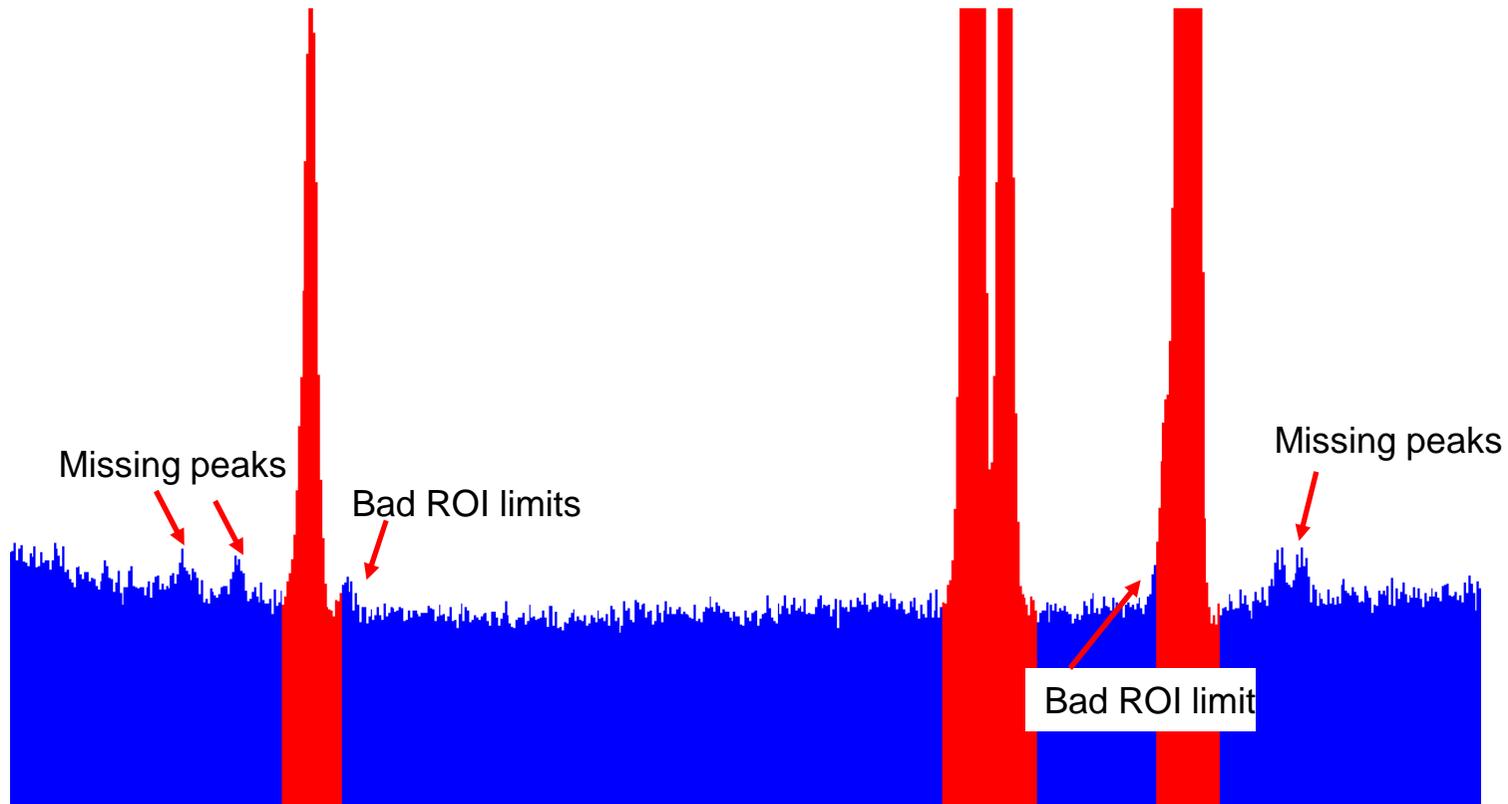
First derivative : change of sign at the maximum peak position

Second derivative : negative minimum at the position peak



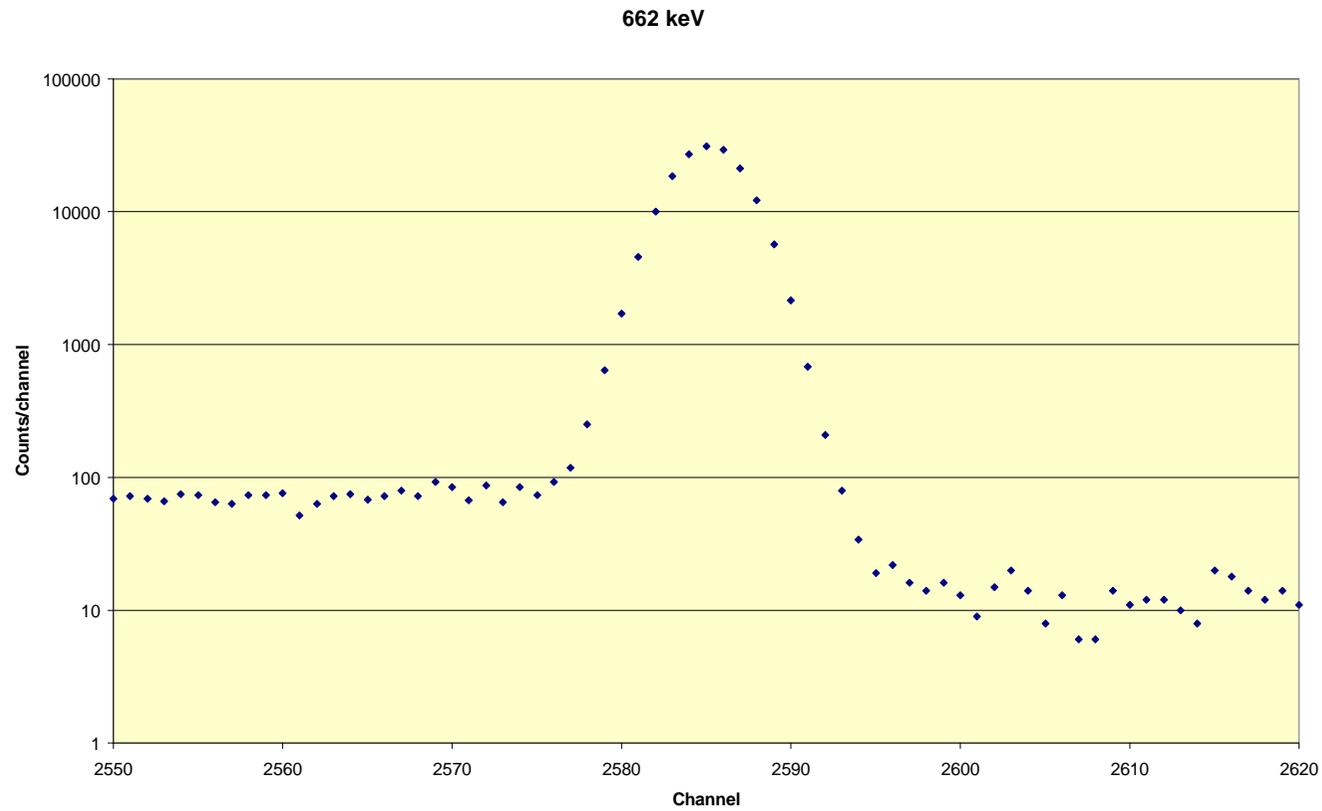
AUTOMATIC PEAK SEARCH

Result of automatic peak search



Expert visual checking recommended !

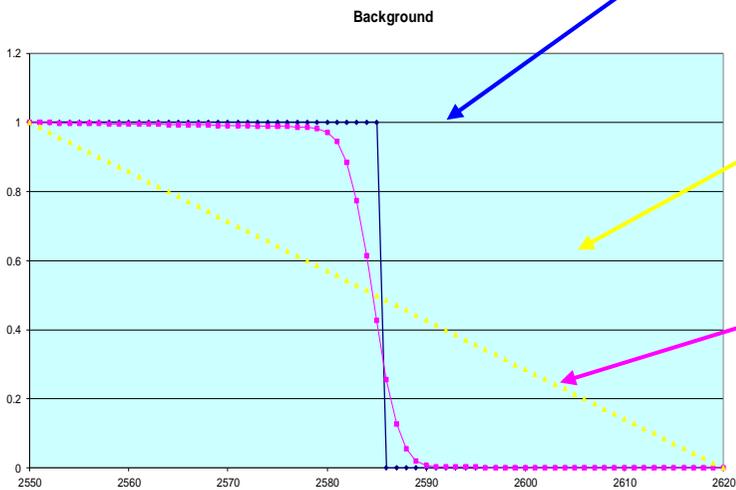
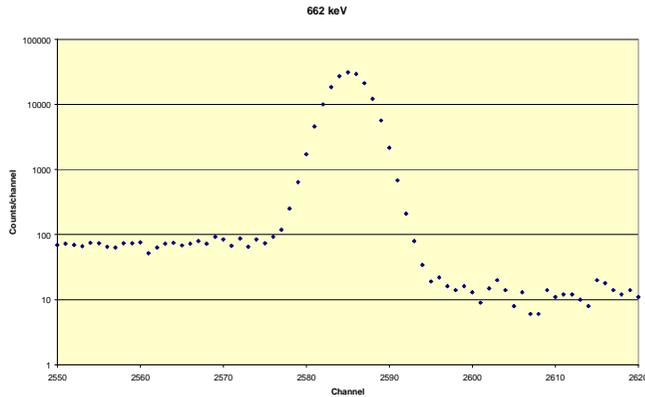
Processing of a single peak ROI



- Peak superimposed on a continuous background

Background

Peak with energy E_0 in region $[f, l]$ first channel = f ; last channel = l



Step :

$$F(i) = 1 \text{ for } E < E_0$$

$$F(i) = 0 \text{ for } E > E_0$$

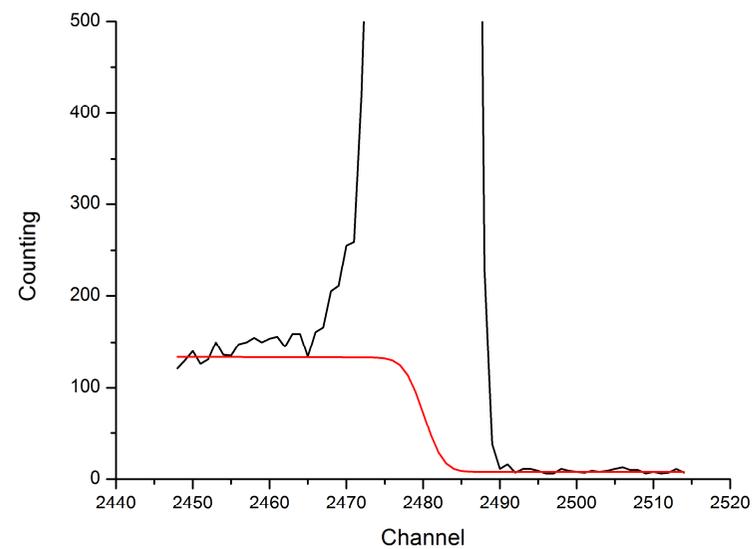
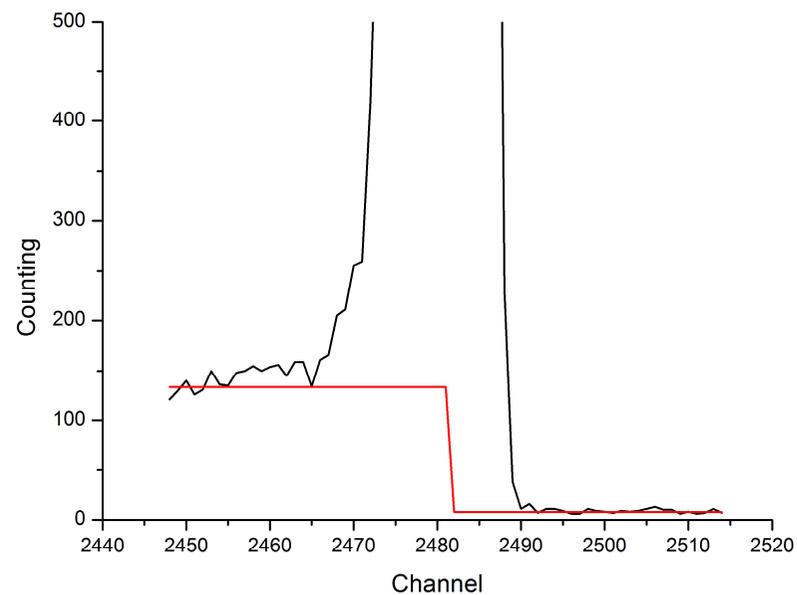
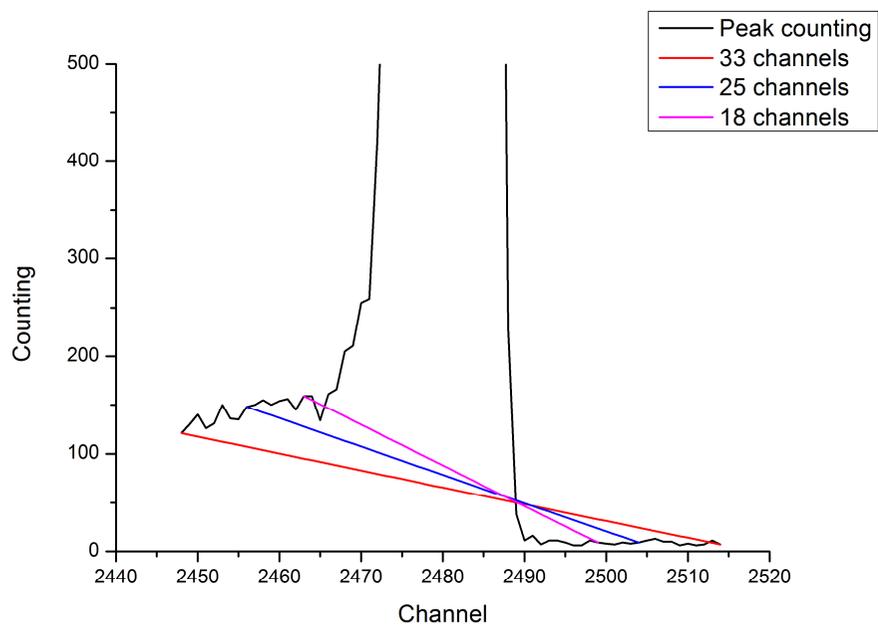
Linear :

$$F(i) = 1 - \frac{(i - f)}{(l - f)}$$

Galton curve :

$$F(i) = 1 - \frac{\sum_{j=f}^i Y_j}{\sum_{j=f}^l Y_j}$$

Background



Background type	Gross area	Net area
Linear (33 channels)	943188	938907
Linear (25 channels)	942037	938191
Linear (18 channels)	940938	937830
Step (33 channels)	943188	938385
Galton (33 channels)	943188	938577

General characteristics

Energy resolution

Peak area

Attenuation

Corrective factors

Use of screen to reduce counting rate (X-rays, Bremsstrahlung)

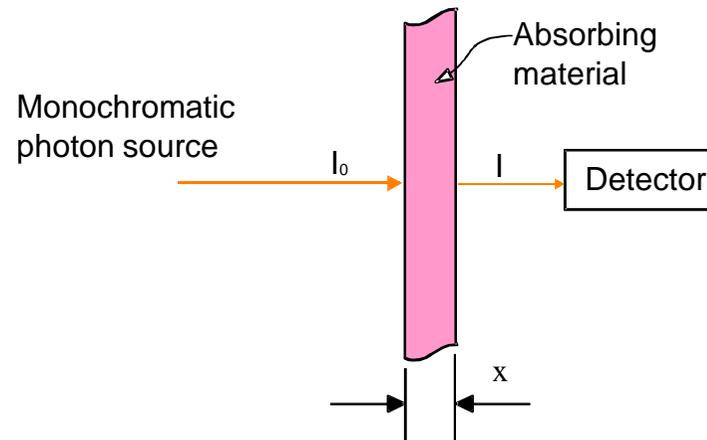
Beer-Lambert law

Mass attenuation coefficients

Beer-Lambert law

Attenuation of a narrow parallel photon beam

$$I(x) = I_0 \cdot e^{-\mu x} = I_0 \cdot e^{-\frac{\mu}{\rho} \rho x}$$



ρ = density ($g.cm^{-3}$)

μ = linear attenuation coefficient of material i for energy E (cm^{-1})

ρx = mass thickness ($g.cm^{-2}$)

μ / ρ = mass attenuation coefficient ($cm^2.g^{-1}$)

μ depends on **E** and **Z**

Attenuation coefficients

Practical parameter : Attenuation (absorption + scattering) coefficient

Partial interaction coefficients depend on the energy, E , and the material (Z):

Photoelectric absorption: $\tau_i(E)$ $\tau \approx \text{const} \cdot Z^{4.5} \cdot E^{-3}$ (dominant at low energies)

Compton scattering: $\sigma_i(E)$ $\sigma \approx \text{const} \cdot Z \cdot E^{-1}$

Pair production effect: $\kappa_i(E)$ $\kappa \approx \text{const} \cdot Z^2$ (only if $E > 1022 \text{ keV}$)

For practical use : tables function of Z and E

Tables : cross sections ($1 \text{ barn} = 10^{-24} \text{ cm}^2$) or mass attenuation ($\text{cm}^2 \cdot \text{g}^{-1}$)

Interaction	Linear attenuation coefficient (cm^{-1})	Mass attenuation coefficient ($\text{cm}^2 \cdot \text{g}^{-1}$)
Photoelectric	τ	τ/ρ
Compton	σ	σ/ρ
Pair production	κ	κ/ρ
Total	$\mu = \tau + \sigma + \kappa$	$\mu/\rho = \tau/\rho + \sigma/\rho + \kappa/\rho$

Attenuation coefficients

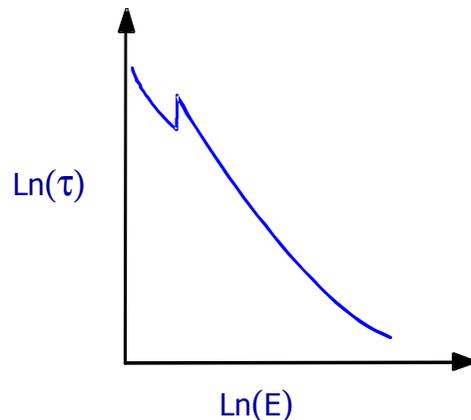
Photoelectric absorption coefficient = sum of photoelectric effect in each electronic shell (subshells):

$$\tau = \tau_K + (\tau_{L1} + \tau_{L2} + \tau_{L3}) + (\tau_{M1} + \tau_{M2} + \tau_{M3} + \tau_{M4} + \tau_{M5}) + \dots$$

If $E < \text{binding energy of shell } i$, $\tau_i = 0$

For $E = E_i$: absorption discontinuity: maximum ionisation probability in shell i

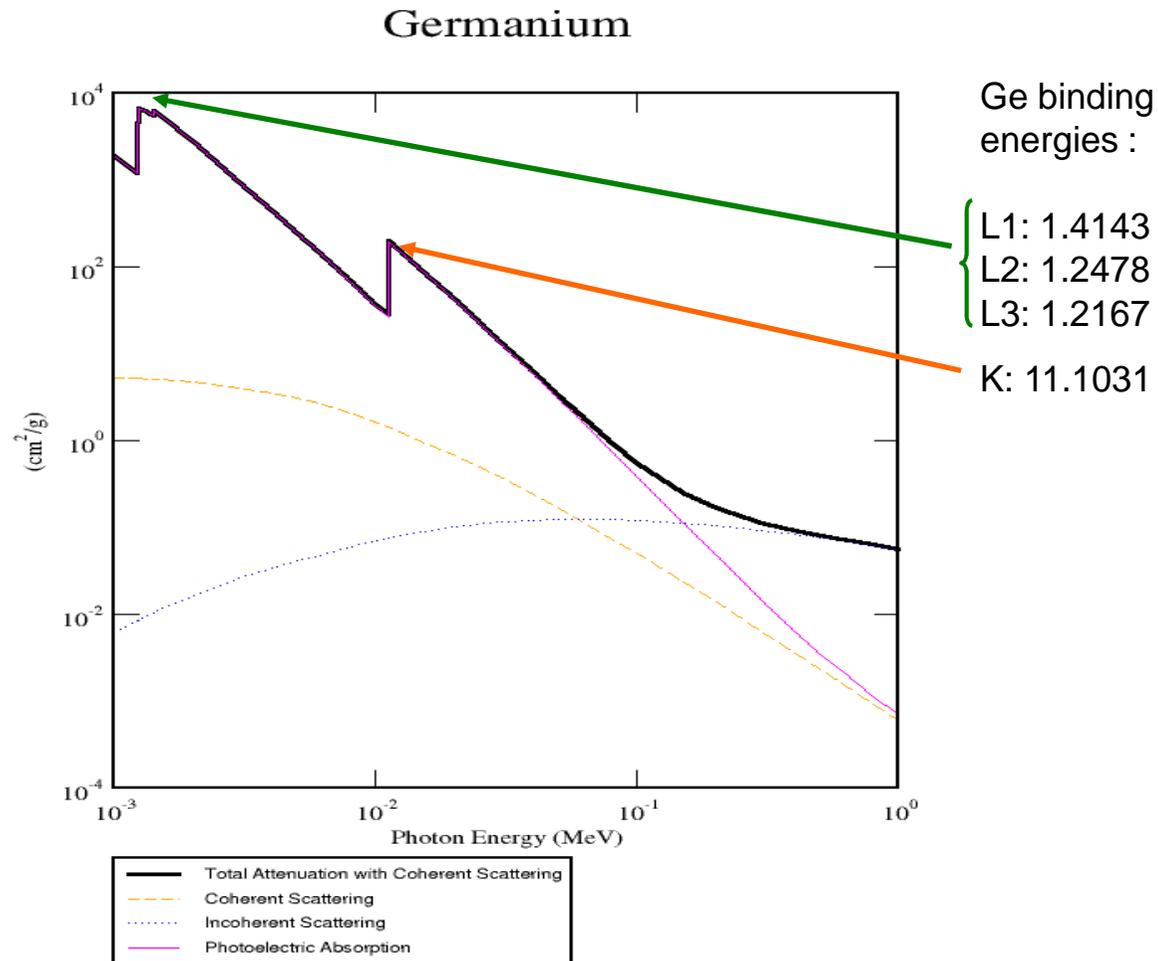
τ variation versus the energy shows discontinuities corresponding to binding energies of electrons shells and subshells K, L, M...



Since $\mu = \tau + \sigma + \kappa$
 μ has the same discontinuities, function of the material atomic structure (Z)

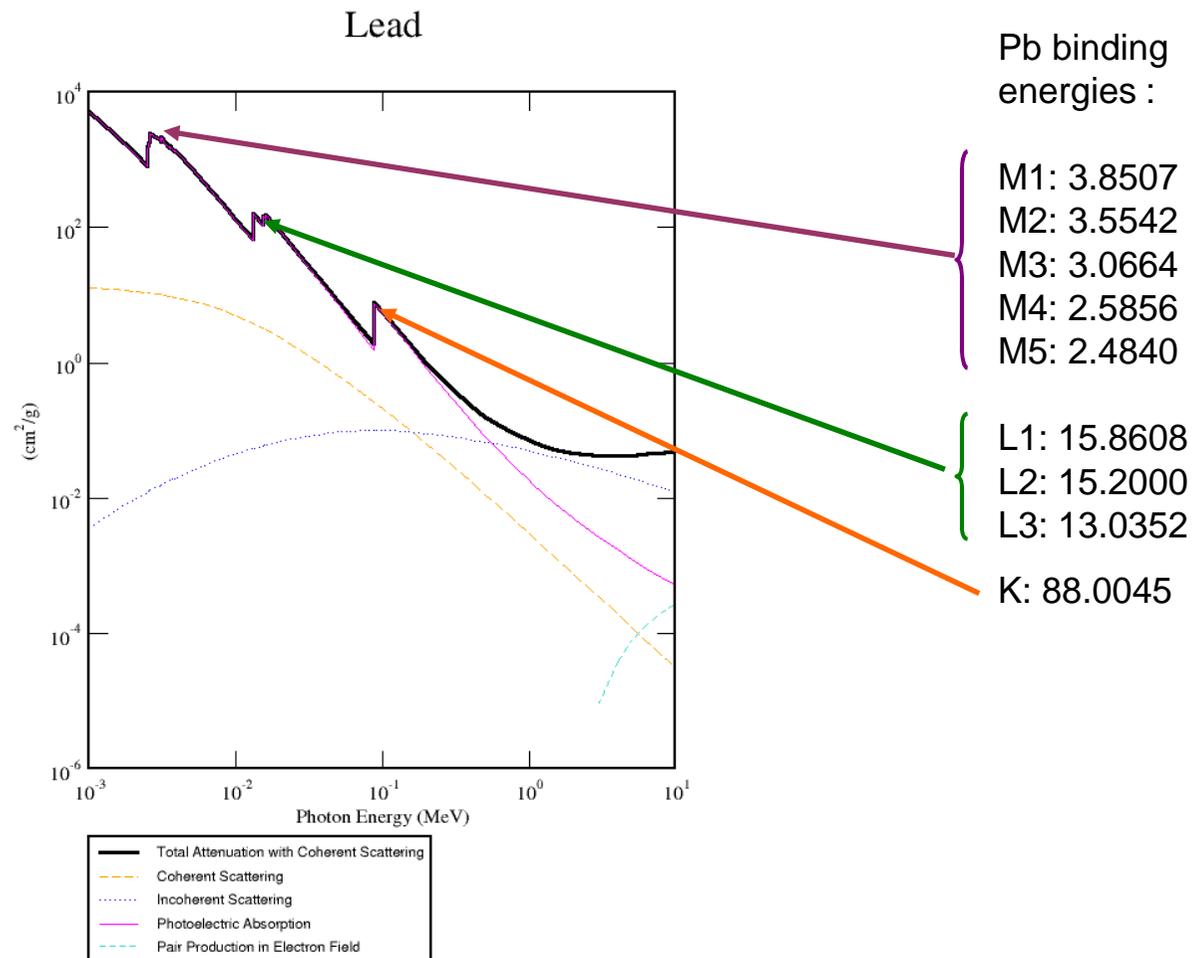
Attenuation coefficients

Germanium mass attenuation coefficient



Attenuation coefficients

Lead mass attenuation coefficient



Composition known -> calculation

Composition unknown -> measurement

Calculation: Attenuation coefficients table

- XCOM (NIST Database)
- Example for HCl 1N

Attenuation coefficients

Defining the mass fraction of each compound for HCl 1N:

Matrix : HCl 1N = 1 mole of HCl in 1 liter of solution

HCl 1N density = 1.016 (1L = 1016 g)

Mass of one HCl mole = $1 + 35.45 = 36.45$ g

Resulting input parameters for XCOM

Compound 1: HCl

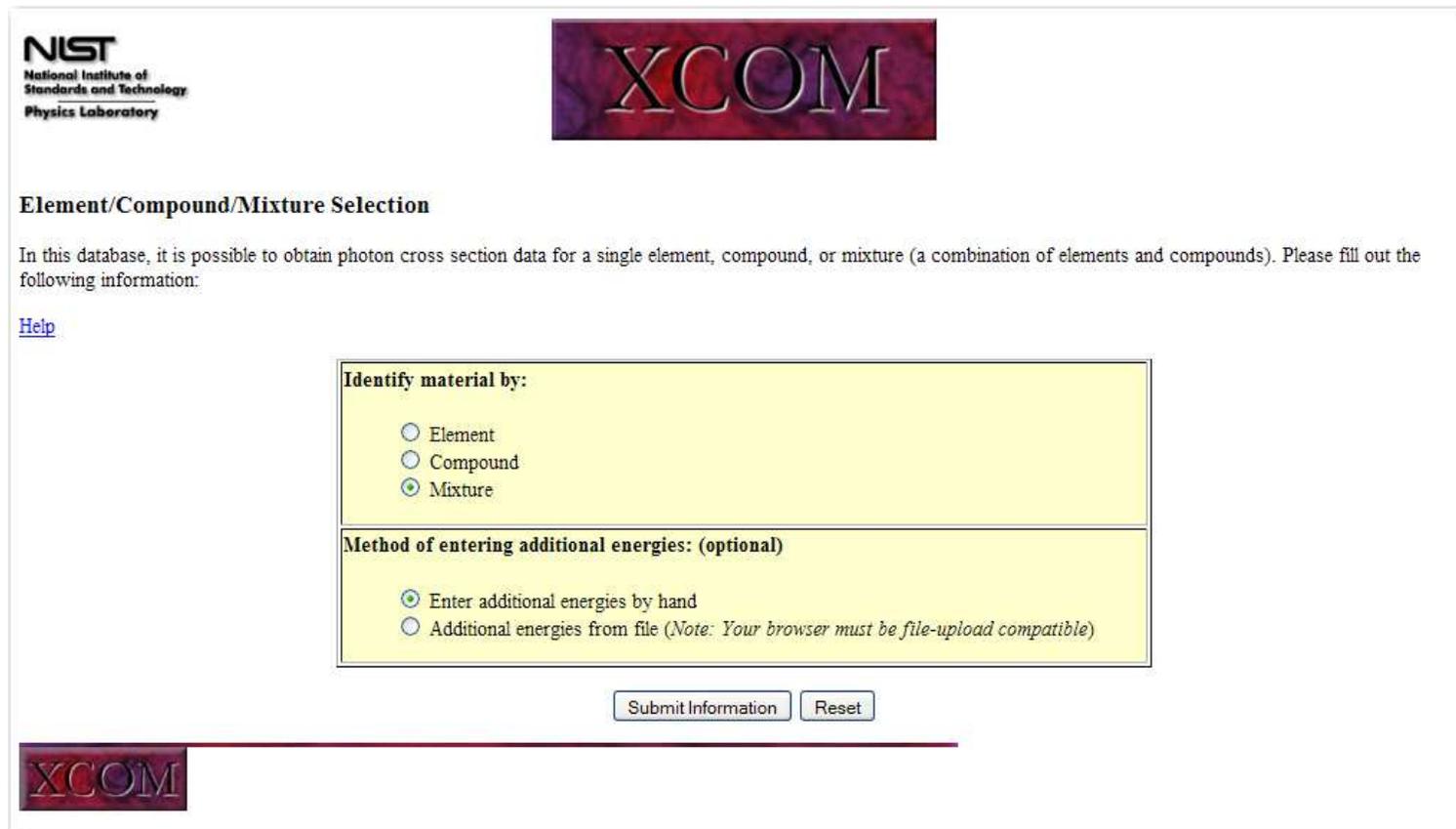
Mass fraction: 36.45

Compound 2: H₂O

Mass fraction = $1016 - 36.45 = 979.55$

Attenuation coefficients

<https://physics.nist.gov/PhysRefData/Xcom/html/xcom1.html>



The screenshot shows the NIST XCOM website interface. At the top left is the NIST logo (National Institute of Standards and Technology, Physics Laboratory). At the top center is a large 'XCOM' logo. Below the logo is the section 'Element/Compound/Mixture Selection'. A paragraph explains that the database provides photon cross section data for single elements, compounds, or mixtures, and asks the user to fill out the following information. A blue 'Help' link is provided. The main form area is divided into two sections: 'Identify material by:' with radio buttons for 'Element', 'Compound', and 'Mixture' (selected); and 'Method of entering additional energies: (optional)' with radio buttons for 'Enter additional energies by hand' (selected) and 'Additional energies from file (Note: Your browser must be file-upload compatible)'. At the bottom of the form are 'Submit Information' and 'Reset' buttons. A small 'XCOM' logo is also visible at the bottom left of the page.

NIST
National Institute of
Standards and Technology
Physics Laboratory

XCOM

Element/Compound/Mixture Selection

In this database, it is possible to obtain photon cross section data for a single element, compound, or mixture (a combination of elements and compounds). Please fill out the following information:

[Help](#)

Identify material by:

- Element
- Compound
- Mixture

Method of entering additional energies: (optional)

- Enter additional energies by hand
- Additional energies from file (*Note: Your browser must be file-upload compatible*)

XCOM

Attenuation coefficients

Fill out the form to select the data to be displayed:

[Help](#)

Enter the formulae and relative weights separated by a space for each compound. One compound per line. For example:
 ^{220}Rn 0.9
 ^{226}Ra 0.1

Note: Weights not summing to 1 will be normalized.

Optional output title:

Graph options:

- Total Attenuation with Coherent Scattering
- Total Attenuation without Coherent Scattering
- Coherent Scattering
- Incoherent Scattering
- Photoelectric Absorption
- Pair Production in Nuclear Field
- Pair Production in Electron Field
- None

Additional energies in MeV: (optional) (up to 75 allowed)

Note: Energies must be between 0.001 - 100000 MeV (1 keV - 100 GeV) (only 4 significant figures will be used). One energy per line. Blank lines will be ignored.

Include the standard grid

Energy Range:

Minimum: MeV
 Maximum: MeV

Constituents (Atomic Number : Fraction by Weight)

Z=1 : 0.108875
 Z=8 : 0.856240
 Z=17 : 0.034884

To download data in spreadsheet (array) form, choose a delimiter and use the checkboxes in the table heading. After down

Delimiter:

- space
- | (vertical bar)
- tab
- newline

Edge	(required) Photon Energy MeV	Scattering		Photoelectric Absorption cm ² /g	Pair Production		Total Attenuation	
		<input type="checkbox"/>	<input type="checkbox"/>		<input type="checkbox"/>	<input type="checkbox"/>	<input type="checkbox"/>	<input type="checkbox"/>
		Coherent cm ² /g	Incoherent cm ² /g		In Nuclear Field cm ² /g	In Electron Field cm ² /g	With Coherent Scattering cm ² /g	Without Coherent Scattering cm ² /g
	1.000E-02	2.48E-01	1.53E-01	6.73E+00	0.00E+00	0.00E+00	7.14E+00	6.89E+00
	1.500E-02	1.44E-01	1.68E-01	1.92E+00	0.00E+00	0.00E+00	2.24E+00	2.09E+00
	2.000E-02	9.57E-02	1.76E-01	7.80E-01	0.00E+00	0.00E+00	1.05E+00	9.55E-01
	3.000E-02	5.08E-02	1.82E-01	2.15E-01	0.00E+00	0.00E+00	4.47E-01	3.96E-01
	4.000E-02	3.12E-02	1.82E-01	8.52E-02	0.00E+00	0.00E+00	2.98E-01	2.67E-01
	5.000E-02	2.11E-02	1.79E-01	4.14E-02	0.00E+00	0.00E+00	2.42E-01	2.21E-01
	6.000E-02	1.52E-02	1.76E-01	2.29E-02	0.00E+00	0.00E+00	2.14E-01	1.99E-01
	8.000E-02	8.93E-03	1.69E-01	8.99E-03	0.00E+00	0.00E+00	1.87E-01	1.78E-01
	1.000E-01	5.86E-03	1.62E-01	4.35E-03	0.00E+00	0.00E+00	1.72E-01	1.66E-01
	1.220E-01	4.00E-03	1.55E-01	2.28E-03	0.00E+00	0.00E+00	1.61E-01	1.57E-01
	1.500E-01	2.68E-03	1.47E-01	1.17E-03	0.00E+00	0.00E+00	1.51E-01	1.48E-01
	2.000E-01	1.53E-03	1.35E-01	4.67E-04	0.00E+00	0.00E+00	1.37E-01	1.35E-01
	3.000E-01	6.84E-04	1.17E-01	1.34E-04	0.00E+00	0.00E+00	1.18E-01	1.18E-01

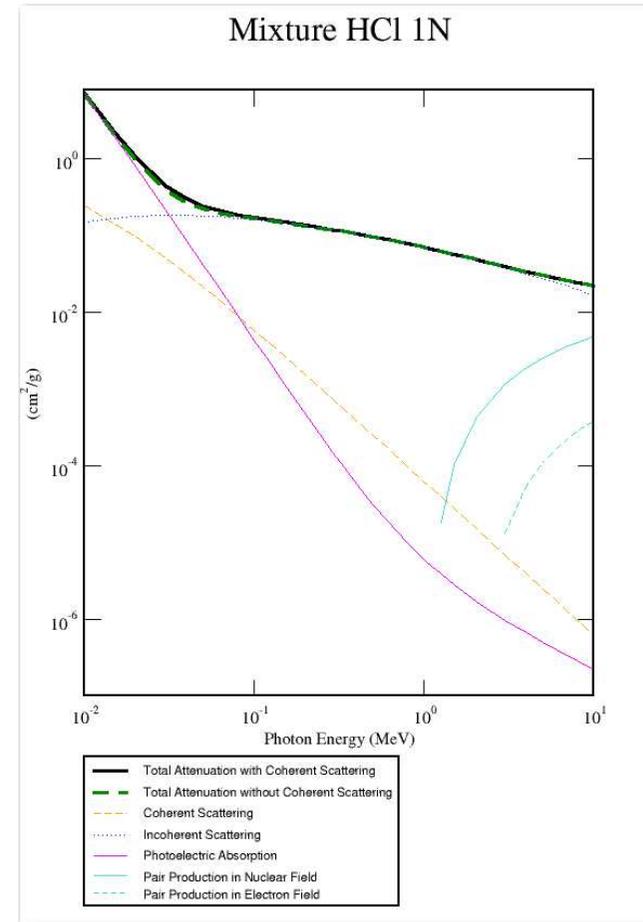
Terminé

Attenuation coefficients

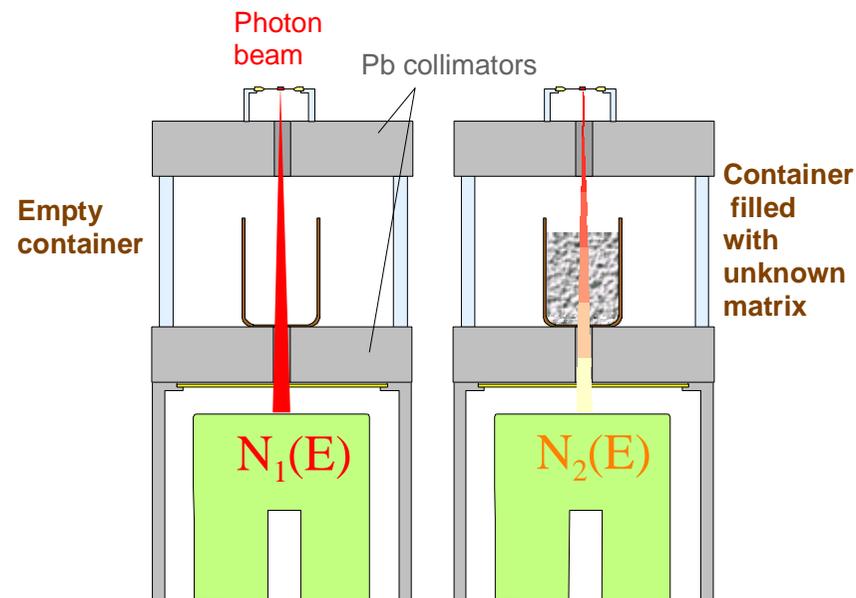
XCOM results

tab
newline
Download data Reset

Edge	(required) Photon Energy MeV	Scattering		Photoelectric Absorption cm ² /g	Pair Production		Total Attenuation	
		Coherent cm ² /g	Incoherent cm ² /g		In Nuclear Field cm ² /g	In Electron Field cm ² /g	With Coherent Scattering cm ² /g	Without Coherent Scattering cm ² /g
		cm ² /g	cm ² /g		cm ² /g	cm ² /g	cm ² /g	cm ² /g
	1.000E-03	1.43E+00	1.31E-02	4.03E+03	0.00E+00	0.00E+00	4.03E+03	4.03E+03
	1.500E-03	1.32E+00	2.65E-02	1.36E+03	0.00E+00	0.00E+00	1.36E+03	1.36E+03
	2.000E-03	1.20E+00	4.14E-02	6.10E+02	0.00E+00	0.00E+00	6.11E+02	6.10E+02
	2.822E-03	9.89E-01	6.52E-02	2.27E+02	0.00E+00	0.00E+00	2.28E+02	2.27E+02
17 K	2.822E-03	9.89E-01	6.52E-02	2.78E+02	0.00E+00	0.00E+00	2.79E+02	2.78E+02
	3.000E-03	9.47E-01	6.99E-02	2.36E+02	0.00E+00	0.00E+00	2.37E+02	2.36E+02
	4.000E-03	7.41E-01	9.32E-02	1.04E+02	0.00E+00	0.00E+00	1.04E+02	1.04E+02
	5.000E-03	5.86E-01	1.11E-01	5.40E+01	0.00E+00	0.00E+00	5.47E+01	5.41E+01
	6.000E-03	4.75E-01	1.24E-01	3.15E+01	0.00E+00	0.00E+00	3.21E+01	3.16E+01
	8.000E-03	3.31E-01	1.42E-01	1.33E+01	0.00E+00	0.00E+00	1.38E+01	1.34E+01
	1.000E-02	2.48E-01	1.55E-01	6.73E+00	0.00E+00	0.00E+00	7.14E+00	6.89E+00
	1.500E-02	1.44E-01	1.68E-01	1.92E+00	0.00E+00	0.00E+00	2.24E+00	2.09E+00
	2.000E-02	9.57E-02	1.76E-01	7.80E-01	0.00E+00	0.00E+00	1.05E+00	9.55E-01
	3.000E-02	5.08E-02	1.82E-01	2.15E-01	0.00E+00	0.00E+00	4.47E-01	3.96E-01
	4.000E-02	3.12E-02	1.82E-01	8.52E-02	0.00E+00	0.00E+00	2.98E-01	2.67E-01
	5.000E-02	2.11E-02	1.79E-01	4.14E-02	0.00E+00	0.00E+00	2.42E-01	2.21E-01
	6.000E-02	1.52E-02	1.76E-01	2.29E-02	0.00E+00	0.00E+00	2.14E-01	1.99E-01
	8.000E-02	8.93E-03	1.69E-01	8.99E-03	0.00E+00	0.00E+00	1.87E-01	1.78E-01
	1.000E-01	5.86E-03	1.62E-01	4.35E-03	0.00E+00	0.00E+00	1.72E-01	1.66E-01
	1.500E-01	2.68E-03	1.47E-01	1.17E-03	0.00E+00	0.00E+00	1.51E-01	1.48E-01
	2.000E-01	1.53E-03	1.35E-01	4.67E-04	0.00E+00	0.00E+00	1.37E-01	1.35E-01
	3.000E-01	6.84E-04	1.17E-01	1.34E-04	0.00E+00	0.00E+00	1.18E-01	1.18E-01
	4.000E-01	3.86E-04	1.05E-01	5.76E-05	0.00E+00	0.00E+00	1.06E-01	1.05E-01
	5.000E-01	2.48E-04	9.62E-02	3.12E-05	0.00E+00	0.00E+00	9.65E-02	9.63E-02
	5.140E-01	2.34E-04	9.51E-02	2.90E-05	0.00E+00	0.00E+00	9.54E-02	9.52E-02
	6.000E-01	1.72E-04	8.90E-02	1.95E-05	0.00E+00	0.00E+00	8.92E-02	8.91E-02
	8.000E-01	9.69E-05	7.82E-02	9.83E-06	0.00E+00	0.00E+00	7.84E-02	7.83E-02
	1.000E+00	6.20E-05	7.04E-02	6.11E-06	0.00E+00	0.00E+00	7.04E-02	7.04E-02
	1.022E+00	5.94E-05	6.96E-02	5.73E-06	0.00E+00	0.00E+00	6.97E-02	6.96E-02
	1.173E+00	4.51E-05	6.50E-02	4.30E-06	6.51E-06	0.00E+00	6.51E-02	6.50E-02
	1.250E+00	3.97E-05	6.29E-02	3.89E-06	1.86E-05	0.00E+00	6.30E-02	6.29E-02
	1.331E+00	3.50E-05	6.09E-02	3.50E-06	3.87E-05	0.00E+00	6.10E-02	6.10E-02
	1.500E+00	2.76E-05	5.73E-02	2.83E-06	1.03E-04	0.00E+00	5.73E-02	5.73E-02



Principle : Use the unknown matrix and a collimated photon beam



Two successive measurements

- Empty container
- Container filled with unknown matrix with thickness x

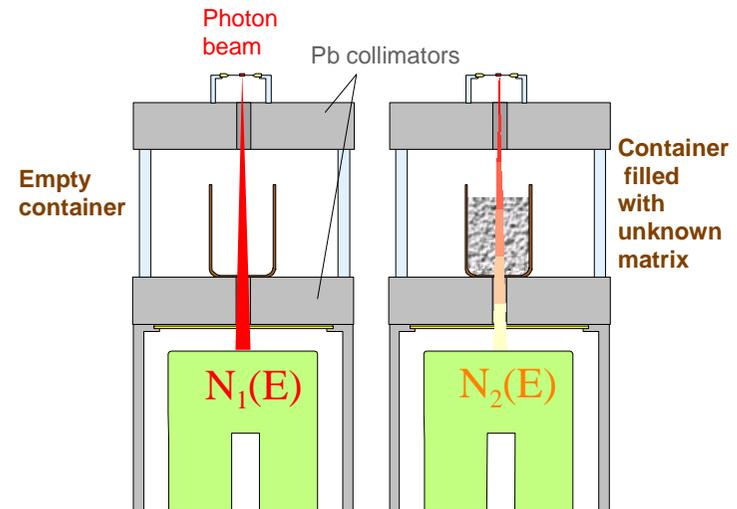
Experimental measurement

For each energy:

$$N_2(E) = N_1(E) \cdot \exp(-\mu(E) \cdot x)$$

Thus:

$$\mu(E) = \frac{1}{x} \ln \left(\frac{N_2(E)}{N_1(E)} \right)$$

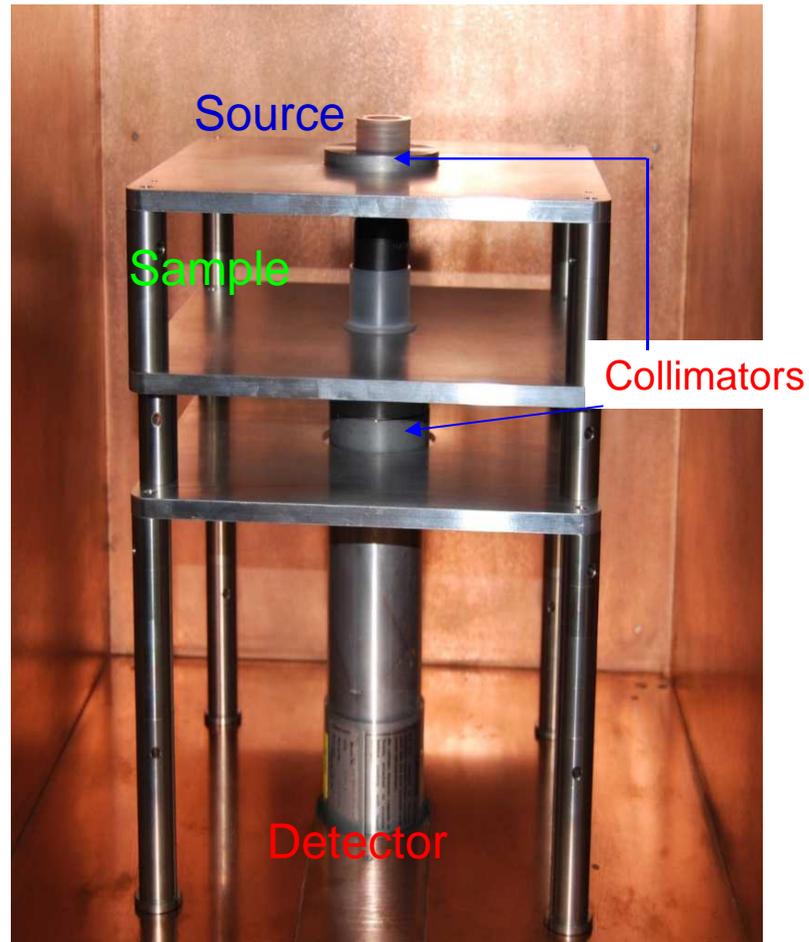


Associated relative uncertainty:

$$\frac{u^2(\mu)}{\mu^2} = \frac{u^2(x)}{x^2} + \frac{1}{\ln^2 \left(\frac{N_0(E)}{N(E)} \right)} \cdot \left(\frac{u^2(N_0(E))}{N_0^2(E)} + \frac{u^2(N(E))}{N^2(E)} \right)$$

Experimental measurement

Dedicated experimental arrangement :



Example of steel sample

⁶⁰Co in steel sample

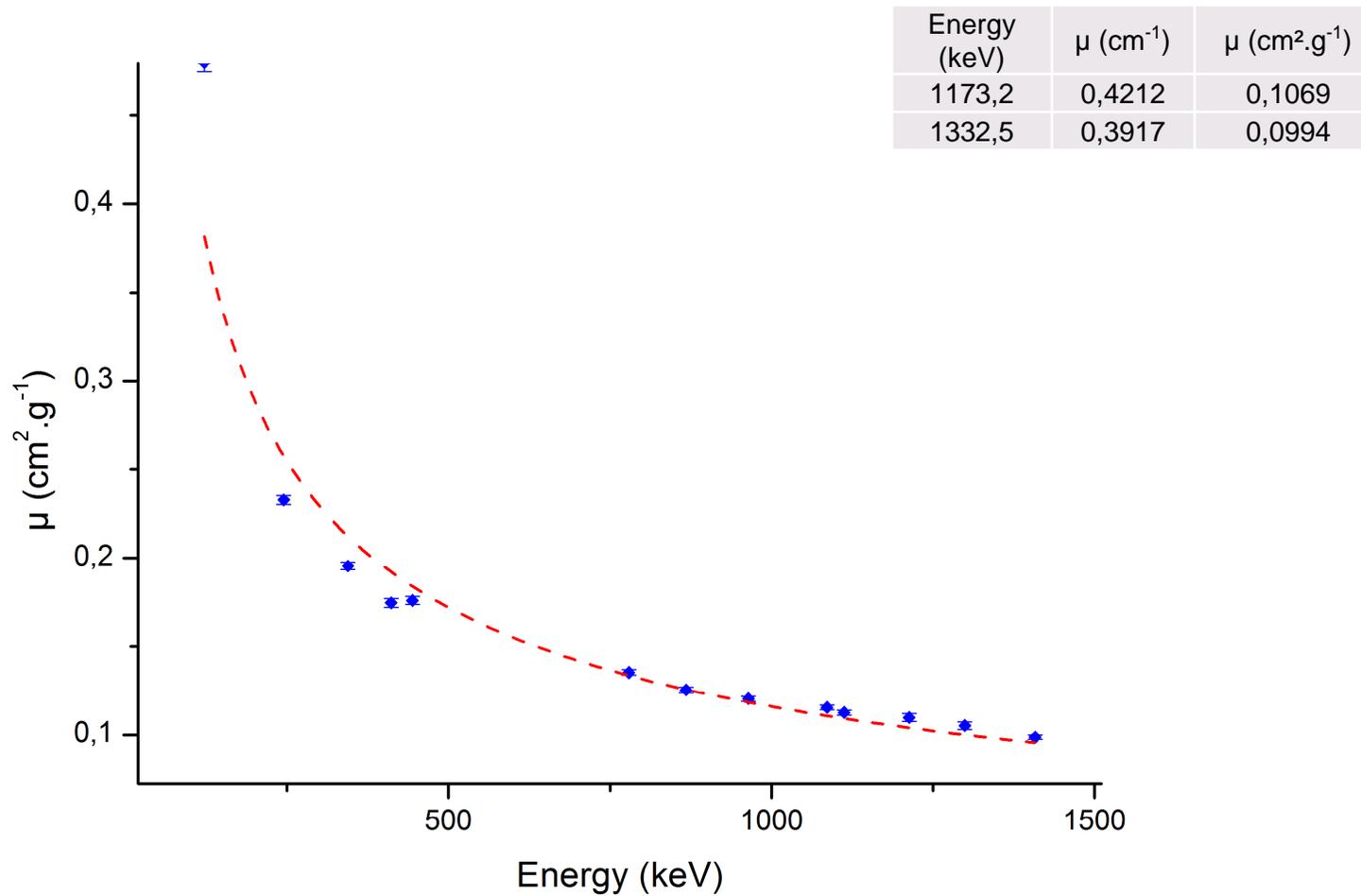
Eu-152		13,522	0,016								
		4,2671E+08	5,0E+05			1 a =	j				
		7,1119E+06	8,4E+03			1 a =	s				
		1,1853E+05	1,4E+02								
		4,9388E+03	5,8E+00								
		1,3522E+01	1,6E-02								
				Date et heure de référence		25/03/2018 12:00					
Eu-152 without material			Eu-152 with material			Attenuation results					
						Material mass (m)		75,8032	g		
						Stand. unc.		0,0008	g		
Start	26/03/2018		Start	27/03/2018		Diameter (D)		3,5	cm		
Real time	80226,6		Real time	77741,46		Stand. unc.		0,02	cm		
Active time	77393,3		Active time	74906,12		Height (x)		1	cm		
% TM	3,5%		% TM	3,6%		Stand. unc.		0,01	cm		
Half-life corr.	1,0001E+00		Half-life corr.	1,0003E+00		Density (ρ)		3,94	g.cm ⁻³		
Decay corr.	1,0001E+00		Decay corr.	1,0001E+00		Stand. unc.		0,01	g.cm ⁻³		
Energy (keV)	NET	Rel. Stand. Unc	NET	Rel. Stand. Unc	I/I0	el. Stand. Un	μ (cm ⁻¹)	el. Stand. Un	μ (cm ² .g ⁻¹)	Rel. Stand. Unc	
121,78	2248977	0,1%	328827	0,2%	0,1511	0,2%	1,890	1,0%	0,480	1,0%	
244,70	412482	0,2%	159477	0,3%	0,3995	0,3%	0,917	1,1%	0,233	1,1%	
344,28	1096300	0,1%	491001	0,1%	0,4628	0,2%	0,770	1,0%	0,196	1,1%	
411,12	66423	0,4%	32317	0,6%	0,5028	0,7%	0,688	1,4%	0,175	1,4%	
443,97	108192	0,3%	52366	0,4%	0,5001	0,5%	0,693	1,3%	0,176	1,3%	
778,90	399743	0,2%	226912	0,2%	0,5866	0,3%	0,533	1,1%	0,135	1,2%	
867,38	131623	0,3%	77692	0,4%	0,6099	0,5%	0,494	1,4%	0,125	1,4%	
964,08	454460	0,1%	273614	0,2%	0,6221	0,2%	0,475	1,1%	0,120	1,2%	
1085,84	371627	0,2%	228130	0,2%	0,6343	0,3%	0,455	1,2%	0,116	1,2%	
1112,08	428505	0,2%	266091	0,2%	0,6417	0,2%	0,444	1,1%	0,113	1,2%	
1212,95	43790	0,5%	27500	0,6%	0,6489	0,8%	0,432	2,0%	0,110	2,1%	
1299,14	53310	0,4%	34077	0,5%	0,6605	0,7%	0,415	1,9%	0,105	2,0%	
1408,01	638601	0,1%	418820	0,2%	0,6777	0,2%	0,389	1,1%	0,099	1,2%	

$$\mu(\text{cm}^{-1}) = \frac{-\ln\left(\frac{I}{I_0}\right)}{x}$$

$$\mu(\text{cm}^1 \cdot \text{g}^{-1}) = \frac{\mu(\text{cm}^{-1})}{\rho}$$

Example of steel sample

Example of steel sample



General characteristics

Energy resolution

Peak area

Attenuation

Corrective factors

Half-life correction

Decay correction (short half-lives)

Attenuation correction (screen)

Self-attenuation correction (see “Self attenuation” section)

Geometry correction (source-detector distance)

Coincidence summing correction (see “Coincidence summing” section)

Background peaks

Escape peaks (annihilation or characteristic Ge X-rays)

Dead time correction

In-flight annihilation correction (Beta plus emitters)

Measurement at a date different from the reference time
To get the counting rate at the reference time: C_T

$T_{1/2}$: radionuclide half-life

T_0 : reference time

T_m : acquisition start time

$$C_T = \exp\left(\ln(2) \cdot \frac{(T_m - T_0)}{T_{1/2}}\right)$$

The radionuclide decays during the measurement

If T_r not negligible versus $T_{1/2}$:

$$C_D = \frac{\ln(2) \cdot (T_r/T_{1/2})}{1 - \exp(-\ln(2) \cdot (T_r/T_{1/2}))}$$

Or: $C_D = \frac{\lambda \cdot T_r}{1 - \exp(-\lambda \cdot T_r)}$ where $\lambda = \frac{\ln(2)}{T_{1/2}}$

$T_{1/2}$: radionuclide half-life

T_r : acquisition time (clock - real time !)

T_m : acquisition start time

Efficiency calibration

- Characteristic of the detector.
- Defined for an incident energy and source-detector **geometry**



Sample with different geometry to be measured

- - **self-absorption**
- - geometry variation (solid angle)

Calibration sources



Calibration sources

Large number of geometries and matrices :

- point source

- disk-shape samples (filters)

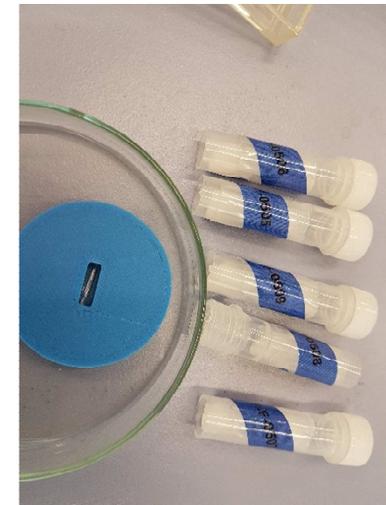


- cylindrical geometries

- liquid

- solid

- gas



- Marinelli geometries (Environment)

- liquid

- solid

- gas



Source-to-detector distance can be adapted according to the counting rate

Small distance : Higher count rate

- ❑ -> Smaller counting time
- ❑ -> Less influence of the environmental background
- ❑ -> Possibility to use smaller samples (less self-absorption)

Large distance : Lower count rate

- ❑ -> Less pile-up
- ❑ -> Reduction of the coincidence summing effect
- ❑ -> Accuracy of the position (relative uncertainty)

Compromise between counting time and corrective factors

Calibration geometry :

Full-energy peak efficiency : $\varepsilon_0(E) = \varepsilon_0^G \cdot \varepsilon^I(E)$

Where $\varepsilon_0^G = \frac{\Omega_0}{4\pi}$

Measurement geometry :

Full-energy peak efficiency : $\varepsilon_m(E) = \varepsilon_m^G \cdot \varepsilon^I(E)$

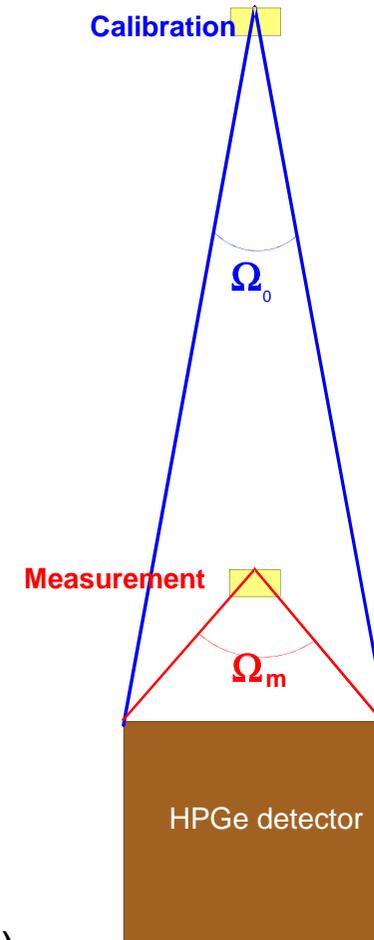
Where $\varepsilon_m^G = \frac{\Omega_m}{4\pi}$



$$\varepsilon_m(E) = \varepsilon_0(E) \cdot \frac{\Omega_m}{\Omega_0}$$

Only valid for point sources

Volume: self-attenuation (+ matrix + source dimensions)



Efficiency transfer example

Efficiency calibration at 100 keV = 0.0171

Detector radius : $r = 2$ cm

Source - detector distance for calibration : $d_0 = 10$ cm

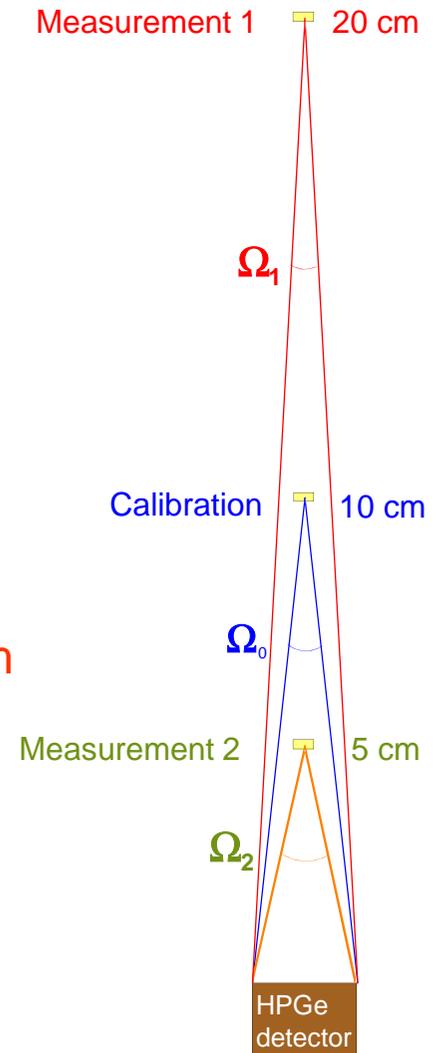
$$\Omega_0 = 2\pi \left(1 - \frac{d_0}{\sqrt{d_0^2 + r^2}} \right) = 0.122$$

Source - detector distance for measurement (m1) : $d_{m1} = 20$ cm

$$\Omega_{m1} = 0.031 \text{ sr} \quad \Omega_{m1}/\Omega_0 = 0.25$$

Source - detector distance for measurement (m2) : $d_{m2} = 5$ cm

$$\Omega_{m2} = 0.449 \text{ sr} \quad \Omega_{m2}/\Omega_0 = 3.7$$



Efficiency transfer example

at 20 cm :

$$\varepsilon_{m1}(100) = \varepsilon_0(100) \cdot \frac{\Omega_{m1}}{\Omega_0} = 0.0171 \cdot \frac{0.031}{0.122} = 0.0043$$

at 5 cm :

$$\varepsilon_{m2}(100) = \varepsilon_0(100) \cdot \frac{\Omega_{m2}}{\Omega_0} = 0.0171 \cdot \frac{0.449}{0.122} = 0.0629$$

Approximation

$$\Omega_0 = 2\pi \left(1 - \frac{d_0}{\sqrt{d_0^2 + r^2}} \right) \approx 2\pi \left(1 - \left(1 - \frac{r^2}{2d^2} \right) \right) = \pi \cdot \frac{r^2}{d^2}$$

$$\frac{\Omega_m}{\Omega_0} = \frac{d_0^2}{d_m^2}$$

$$\varepsilon_{m1}(100) = \varepsilon_0(100) \cdot \frac{d_0^2}{d_{m1}^2} = 0.0171 \cdot \frac{(10)^2}{(20)^2} = 0.0043 \quad \Delta = \text{negligible}$$

$$\varepsilon_{m2}(100) = \varepsilon_0(100) \cdot \frac{d_0^2}{d_{m2}^2} = 0.0171 \cdot \frac{(10)^2}{(5)^2} = 0.0684 \quad \Delta = + 9\%$$

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