

**ICRM GSWG**



Meeting of the ICRM Gamma Spectrometry Working Group  
Monte Carlo benchmark on coincidence summing corrections  
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## **Contribution of Uncertainty of Decay Data and of the Uncertainty of the Dead Layer to the Uncertainty of Coincidence Summing Corrections**

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## Outline

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# 1. Coincidence summing corrections – general considerations

Coincidence summing corrections for a peak of energy  $E_i$  result from:

- *Coincidence losses* due to simultaneous detection of other photons  $E_j$ 
  - Proportional with joint emission probability of the  $E_i$  and  $E_j$  photons  $p_{ij}$  and with the probability of simultaneous detection of  $E_i$  in the peak and of  $E_j$  in the total spectrum
- *Coincidence summing in* in the case when the transition in which  $E_i$  is emitted can also be realized by successive transitions of photons  $E_p$  and  $E_q$ 
  - Proportional with joint emission probability of the  $E_p$  and  $E_q$  photons  $p_{pq}$  and with the probability of simultaneous detection of both  $E_p$  in the peak and of  $E_q$  in the peak

In the case of a point source and negligible angular correlations:

$$F_C(E_i; X) = 1 - \sum_j \frac{p_{ij}}{p_i} \cdot \eta(E_j) + \sum_{j,k} \frac{p_{ijk}}{p_i} \cdot \eta(E_j) \cdot \eta(E_k) - \dots$$

$$+ \sum_{p,q} \frac{p_{pq}}{p_i} \cdot \frac{\varepsilon(E_p) \cdot \varepsilon(E_q)}{\varepsilon(E_i)} - \sum_{p,q,r} \frac{p_{pqr}}{p_i} \cdot \frac{\varepsilon(E_p) \cdot \varepsilon(E_q) \cdot \eta(E_r)}{\varepsilon(E_i)} + \dots$$

$p_i$  – emission probability of the  $E_i$  photon of nuclide  $X$

$p_{ij}, p_{ijk}$  – joint emission probability of photons  $E_i$  and  $E_j$ , respectively  $E_i, E_j, E_k$

$\varepsilon(E)$  – full energy peak efficiency for the photon of energy  $E$

$\eta(E)$  – total efficiency for the photon of energy  $E$

Note:

- In the case of volume sources, the detection efficiencies for group of photons are more involved (angular correlations are neglected):

$$\varepsilon(E_i) \cdot \eta(E_j) \rightarrow \frac{1}{V} \int_V \varepsilon_P(E_i, \vec{r}) \cdot \eta_P(E_j, \vec{r}) dv$$

$$\varepsilon(E_i) \cdot \varepsilon(E_j) \rightarrow \frac{1}{V} \int_V \varepsilon_P(E_i, \vec{r}) \cdot \varepsilon_P(E_j, \vec{r}) dv$$

$\varepsilon_P(E_i, \vec{r})$  and  $\eta_P(E_j, \vec{r})$  = peak and total efficiency for a point source at  $\vec{r}$  (Sima & Arnold, ARI 53 (2000) 51)

## 2. Dependence on decay data parameters

$F_C$  depends on decay scheme parameters through the joint emission probabilities

$p_{ij}$ ,  $p_{ijk}$  depend not only on  $p_i$ ,  $p_j$ ,  $p_k$ , but also on:

- Relative de-excitation of the initial level on other levels
  - The conversion coefficients  $\alpha_K$ ,  $\alpha_L$ ,  $\alpha$
  - The decay branch on each level
  - $\omega_K$  in the case when X-rays contribute (K X rays fluorescence yield)
  - $P_K$  in the case of EC decays (probability of electron capture on K atomic shell)
- ⇒ The explicit dependence of  $p_{ij}$ ,  $p_{ijk}$  on the decay scheme parameters: very complex in the general case
- ⇒ How to evaluate the uncertainty of  $F_C$  due to the uncertainties of the parameters of the decay scheme?
- Uncertainty propagation formula not convenient – nonlinear, correlations
  - Standard Monte Carlo not convenient – decay scheme simulation mixed with radiation transport => the distribution of  $F_C$  values would have inseparable contributions both from the uncertainty of decay parameters and of radiation transport

### 3. Procedure for the evaluation of random decay schemes

Proposed solution (Sima & Lepy, ARI 109 (2016) 493):

- Disentangle (cf. Eq. 1) evaluation of joint emission probabilities from efficiencies (radiation transport)
- Evaluate the required joint detection probabilities (group efficiencies) by a separate, long run Monte Carlo simulation
- Prepare a large set of decay scheme data on the basis of the parameters of the decay scheme and of their uncertainties
  - data source: DDEP ([http://www.nucleide.org/DDEP\\_WG/DDEPdata.htm](http://www.nucleide.org/DDEP_WG/DDEPdata.htm))
  - $p_i, P_K, \omega_K$  – independent Gaussian distributions
  - $\alpha_K, \alpha_L, \alpha$ : two procedures: totally correlated Gaussian, or each extracted independently from a Gaussian distribution
  - Transition probabilities: from  $p_i$  and  $\alpha$ 's
  - Decay branches: probability balance (de-excitation and feeding)
  - Acceptance checks
- For each decay scheme from the set evaluate each of the necessary  $p_i, p_{ij}, p_{ijk}$  – (Sima & Arnold, ARI 66 (2008) 705)
- Using the pre-computed joint detection probabilities, for each decay scheme evaluate  $F_C$  (Eq. 1)
- Analyze the distribution of the set of  $F_C$  values: average, standard deviation

## 4. Results – Contributions of the uncertainty of decay scheme parameters

Computations: all the cases from the  $F_C$  benchmark exercise (2 detectors, 4 geometries, 11 peaks = 208 cases)

Generally very low uncertainties: 151 cases from 208 with uncertainty < 0.1%

Highest  $F_C$  uncertainties: detector B, point source, then detector B, filter

79 keV (B, point): 8.2%

160 keV (B, point): 4.6%

79 keV (B, filter): 4.3%

Uncertainty for  $F_C$  at 79 keV dominated by the uncertainty of  $p_i$  for 79 keV:

2.63 (19) (DDEP)

Uncertainty of  $p_i * F_C$

- This quantity is required for activity computation using the count rate  $N$  ( $N = \epsilon_0 * F_C * p_i * A$ )

Higher uncertainty than for  $F_C$  !

112 from 208 cases with uncertainty < 0.5%

For 80 keV: between 1 and 1.7%

## 5. Conclusions – effects of decay scheme uncertainties

Generally the effect of decay scheme uncertainties on the uncertainty of coincidence summing correction factors  $F_C$  is lower than the contribution of other sources of uncertainty

- In particular cases it is not negligible – low intensity peaks with high coincidence summing from other photons

Realistic evaluation of the contribution of decay scheme uncertainties to  $F_C$ :

- Decouple evaluation of joint emission probabilities from radiation transport
- Prepare sets of random decay schemes based on Monte Carlo simulation of decay scheme parameters
  - In absence of covariance matrices of the parameters, sampling of the parameters should be arranged in line with the procedure of parameters evaluation
  - For each decay scheme evaluate the set of joint emission probabilities
  - Using fixed, independently computed, detection probabilities of groups of photons, evaluate  $F_C$
  - Analyse the statistics of  $F_C$  distributions

## 6. Dead Layer problems in $F_C$ evaluation

At low energy (< 100 keV):

- Dead layer (DL) thickness – essential for peak efficiency (high attenuation of low  $E$  photons in DL)
- Dead layer also important for total efficiency for low  $E$  photons
- Best value of the DL thickness: trial and error based on comparison of peak efficiency from simulations with measured values at low energies

At higher energies:

- Attenuation of photons in the dead layer less important
- Effect of DL on efficiency at higher energies: mainly due to the dependence of the volume of the sensitive region on DL thickness

Additional dependence on DL thickness in the case when coincidence summing effects are important (Arnold & Sima, ARI 60 (2004) 167)

- ⇒ Impact of DL thickness on apparent peak efficiency  $\varepsilon(E)$  for higher energy photons: higher than expected for energy  $E$  in the case when coincidence summing with low energy photons is significant for the peak of energy  $E$ 
  - ⇒ The effect is dominated by the attenuation of the coincident low energy photons

Measurements in PTB (Arnold & Sima, ARI 60 (2004) 167):

- Ba-133 and Eu-152 point sources with two detectors (p and n-type)
- For each detector and source:
  - One measurement with a stainless steel absorber of 0.97 mm
  - One measurement without absorber
- DL thickness evaluated by comparison simulation vs. measurement of peak efficiency at low energy (“conventional dead layer thickness”)

Results:

- p-type detector:
  - Peak shape distorted in the case of measurement without absorber
    - Tail in the high energy part of the peaks affected by coincidences with X-rays (Ba-133 peaks and Eu-152 peaks from EC decay branch)
  - Computed  $F_C$  values for these peaks in disagreement with measured values if the conventional DL thickness is used in calculations
  - No problems in the case of measurements with the absorber
- n-type detector: no problems

Simulations with GEANT3 and PENELOPE do not explain the distortion of the peak shape (Stancu et al., Rom. Rep. Phys. 67 (2015) 465).

# Why $F_C$ and $\eta$ evaluation using conventional dead layer thickness fails?

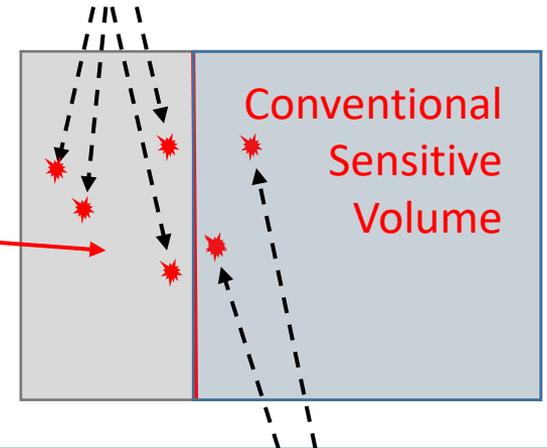
- The sensitive volume for peak efficiency differs from the sensitive volume for total efficiency (and for coincidence losses from a peak)
  - Conventional sensitive volume: all charges produced are collected
  - Sensitive volume for total efficiency: some charge is collected
  - Sensitive volume for total efficiency bigger than sensitive volume for peak efficiency
- ⇒ Dead layer thickness **DLT < DLP**

Difficult to measure DLT

⇒ Usually total efficiency and coincidence losses from peaks are evaluate using DLP (much easier to measure) instead of using DLT

Fraction of the charge produced not collected (no signal in the peak)

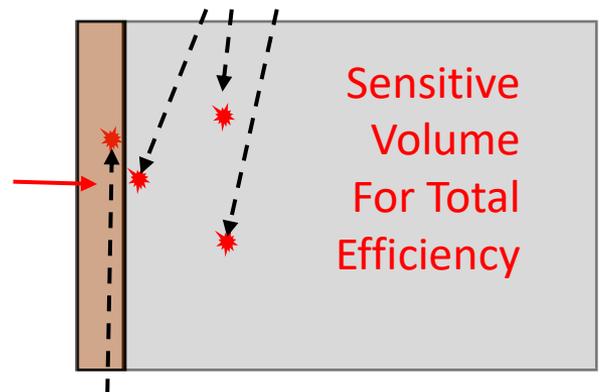
Conventional Dead Layer DLP



Total charge produced contributes to detector signal

Fraction of the charge produced collected (signal in the spectrum)

Dead Layer for the Total Efficiency DLT



No charge produced collected (no signal in the spectrum)

## 7. Consequences of different dead layers for peak and total efficiency

⇒ Effect of  $DLT < DLP$  on the computation of the total efficiency:

- $\eta$  computed using DLP is underestimated – at high energy only due to smaller active volume when using DLP

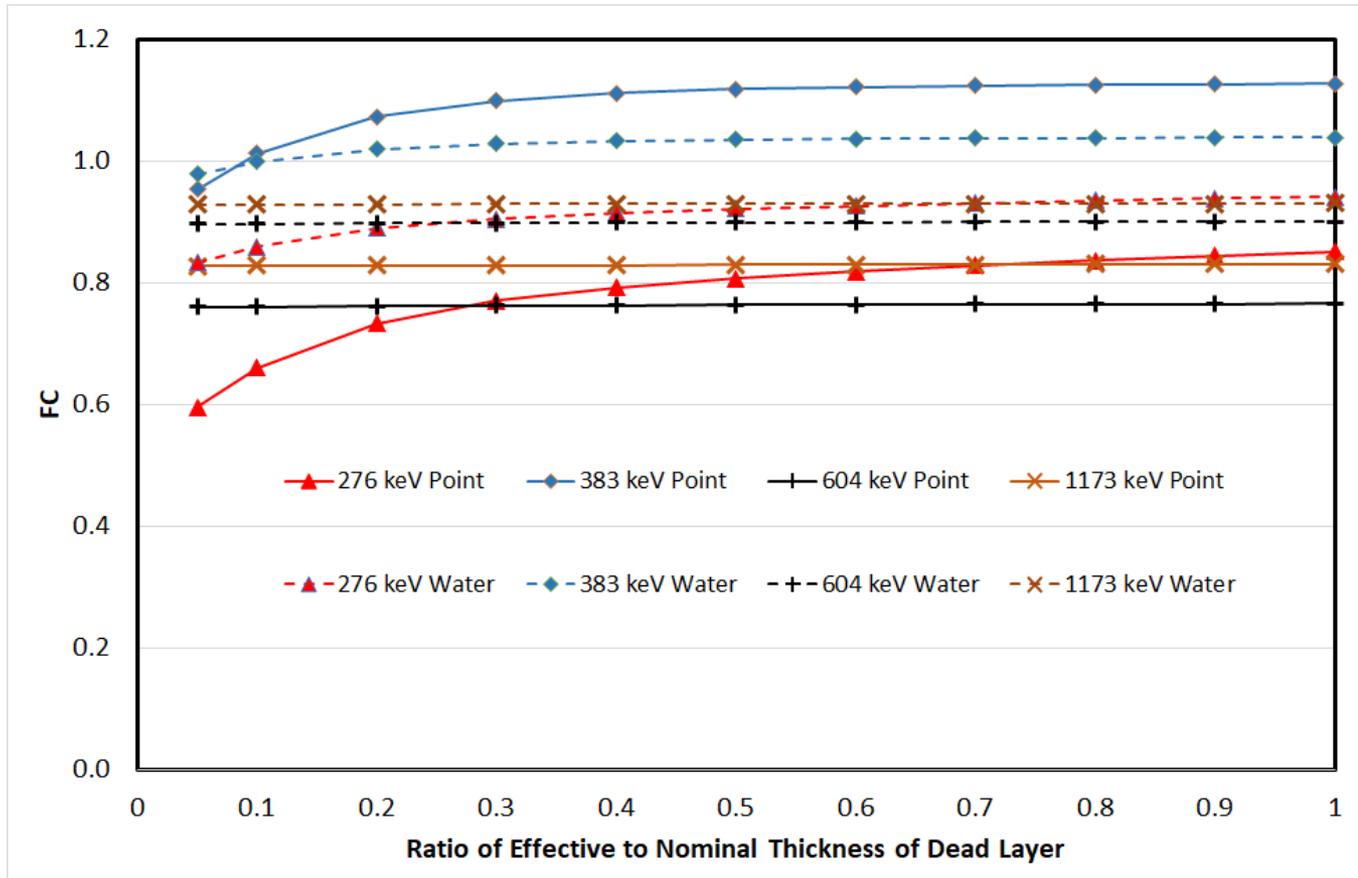
Effect of  $DLT < DLP$  on the computations of  $F_C$ :

- Coincidence summing losses computed using DLP are underestimated
- In the case of high contribution to coincidence losses due to low energy photons, the effect on  $F_C$  is controlled by the dependence on DLT of the sensitive volume for the low energy photons
  - Losses from the peaks of higher energies  $E$  more important than the dependence of  $\eta(E)$  on DLT
- In the case when only high energy photons contribute to coincidence losses, the effect of  $DLT < DLP$  less important

PTB Measurements (Arnold & Sima, ARI 60 (2004) 167):

=> Best results of  $F_C$  evaluation for Ba-133 peaks measured with the p-type detector without absorber:  $DTL/DLP=0.13$

Calculations with GESPECOR for point and water sources (as in Benchmark...)  
 Detector A with nominal dead layer thickness DLP and various effective DLT values



For peaks with low energy photon coincidence summing effects: 15-20 % effect for point source, 5 % for water source

For peaks without low energy photon summing – negligible effect

=> **Good knowledge of the structure of the dead layer required !**

## 8. Conclusions – dead layer problem

- The dead layer has a complex structure
  - ⇒ The sensitive volume for peak efficiency differs from the sensitive volume for total efficiency
  - ⇒ With standard Monte Carlo simulations, in which the signal is completely related to the charge produced in the sensitive volume, it is impossible to reproduce simultaneously the peak and total efficiency, respectively the peak efficiency and the coincidence summing correction factors using a single dead layer.
  - ⇒ The distortion of the peak shape due to coincidences with low energy photons (Arnold & Sima, ARI 60 (2004) 167) cannot be reproduced by standard Monte Carlo simulations (Stancu et al., Rom. Rep. Phys. 67 (2015) 465).
- Realistic solution of the problem: include the simulation of charge collection processes in the software for detector simulation
  - Difficulties: requires knowledge of the electric field distribution, of the impurities distribution, of the charge mobility
- Practical solution: Inclusion of 2 different dead layers in simulations:
  - Conventional dead layer DLP (fixed by peak efficiency measurement at low E)
  - Dead layer for total efficiency DLT, with  $DLT < DLP$
  - DLT may be fixed by comparison of simulations with measurements of  $F_C$  for coincidence losses due to low energy photons