

Laboratoire National Henri Becquerel

LNE-LNHB

Detection efficiency

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Detector efficiency

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 - Definition
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FULL ENERGY PEAK EFFICIENCY

Full-energy Peak Efficiency (FEPE): $\mathcal{E}_{P}(E)$

Ratio of the number of counts in full-energy peak corresponding to energy $E(N_P(E))$, by the number of photons with energy Eemited by the source (F(E))





 $\mathcal{E}_{P}(E)$ depends on the source-detector geometry and on the energy



Geometrical efficiency



 Ω = solid angle between source and detector (*sr*)

For a point source :

$$\Omega = 2\pi \left(1 - \frac{d}{\sqrt{d^2 + r^2}}\right)$$

Ratio of the number of photons emitted towards the detector by the number of photons emited by the source

$$\mathcal{E}_{\rm G} = \frac{\Omega}{4\pi}$$

 ϵ_{G} depends only on the source-detector geometry

Intrinsic efficiency



Difficulty: exact composition badly known

Calculation of the detector FEP efficiency

Transmission probability through material i with thickness x_i : (Beer-Lambert law) :

Thus

Interaction probability in the same material:

$$\implies P_{I}(E, x_{i}) = 1 - P_{T}(E, x_{i})$$

To result in a count in the FEP peak: The photon must :

- be emitted in the Ω solid angle,

- cross the screens (air,n window, dead layer,...) without being absorbed,

- and be totally absorbed in the detector active volume.



Calculation of the detector efficiency



For higher energies:

Total absorption is due to succesive effects : multiple scattering

Thus the calcul is not possible

Calculation of the detector efficiency

Many difficulties for an accurate calculation

- Exact knowledge of the detector parameters: materials (composition) geometry (thickness, position ...)

Data used in the calculation :
 Attenuation coefficients
 Material density

- Semi-conductor effects: parameters and physical interaction :

Electrical field

Electrodes



Radiography of a HPGe detector: Rounded crystal, axially shifted, tilted ...

Experimental FEP efficiency calibration

$$\mathcal{E}_{P}(E) = \frac{N_{P}(E)}{F(E)}$$

$$\mathcal{E}_{P}(E) = \frac{N_{P}(E)}{A \cdot I_{\gamma}(E)}$$

This is performed using standard radionuclides with standardized activity A (Bq) with photon emission intensities, $I\gamma$ well known

 $N_P(E)$: peak net area



 $\epsilon_{P}(E)$ Full-energy peak (FEP) efficiency depends on the energy and on the source-detector geometrical arrrangement

Associated standard uncertainty

Efficiency calibration

$$\mathcal{E}(\mathsf{E}) = \frac{\mathsf{N}(\mathsf{E})}{\mathsf{A} \cdot \mathsf{I}_{\gamma}(\mathsf{E})}$$

$$\left(\frac{\varDelta \varepsilon(\mathsf{E})}{\varepsilon}\right)^2 = \left(\frac{\varDelta \mathsf{N}(\mathsf{E})}{\mathsf{N}}\right)^2 + \left(\frac{\varDelta \mathsf{A}}{\mathsf{A}}\right)^2 + \left(\frac{\varDelta \mathsf{I}(\mathsf{E})}{\mathsf{I}(\mathsf{E})}\right)^2$$

$$\frac{\Delta N(E)}{N(E)} = \frac{\sqrt{N(E)}}{N(E)} = \frac{1}{\sqrt{N(E)}} \qquad \qquad \frac{\Delta A}{A} = 5 \cdot 10^{-3} \qquad \qquad \frac{\Delta I_{\gamma}(E)}{I_{\gamma}(E)} = 1 \cdot 10^{-3}$$

Influence of the peak area :

if N = 10⁴ Δ N/N = 10⁻² -> Δ E/E = 1,1 10⁻²

if N = 10⁵ Δ N/N = 3,1 10⁻³ -> Δ ϵ/ϵ = 6 10⁻³

FEP efficiency calibration

To get an efficiency values at any energy : energy calibration over the whole energy range

1. Use different radionuclides to get energies regularly spaced over the range of interest

Single gamma-ray emitters : ⁵¹Cr (320 keV), ¹³⁷Cs (662 keV) ⁵⁴Mn (834 keV) : one efficiency value per one measurement



Multigamma emitters : ⁶⁰Co, ¹³³Ba, ¹⁵²Eu, ⁵⁶Co : several efficiencies values per one measurement , but coincidence summing effects !

For each energy, discrete values of the FEP efficiency $\varepsilon(E_1)$, $\varepsilon(E_2)$, ... $\varepsilon(E_n)$

2. Computation of the efficiency for

- 1. Local interpolation
- 2. Fitting a mathematical function to the experimental values

FEP efficiency calibration

Efficiency calibration for different source-to-detector distances



Efficiency calibration - Interpolation

Local interpolation

 $E_1 \rightarrow \varepsilon_1$ $E_2 \rightarrow \varepsilon_2$ To get $\varepsilon(F)$ for $F_1 < F_2 < F_3$

Solution 1 : linear interpolation

$$\Rightarrow \varepsilon(E) = \varepsilon_1 + \frac{(E - E_1)}{(E_2 - E_1)} \cdot (\varepsilon_2 - \varepsilon_1)$$

: valid only for close energies

Solution 2 : logarithmic interpolation

$$ln(\varepsilon(E)) = ln(\varepsilon_1) + \frac{(lnE - lnE_1)}{(lnE_2 - lnE_1)} \cdot (ln\varepsilon_2 - ln\varepsilon_1)$$

Efficiency calibration - Interpolation

Local interpolation : example

Determine efficiency for $E = 662 \text{ keV} (^{137}\text{Cs})$ knowing :

1.	$E_1 = 569 \text{ keV} (^{134}\text{Cs})$ $E_2 = 766 \text{ keV} (^{95}\text{Nb})$	$\epsilon_1 = 2.21 \ 10^{-3}$ $\epsilon_2 = 1.67 \ 10^{-3}$
	Linear interpolation : Logarithmic interpolation :	$\epsilon(E) = 1.96 \ 10^{-3}$ $\epsilon(E) = 1.92 \ 10^{-3}$
2.	E ₁ = 122 keV (⁵⁷ Co) E ₂ = 1 173 keV (⁶⁰ Co)	$\epsilon_1 = 8.22 \ 10^{-3} \ \epsilon_2 = 1.13 \ 10^{-3}$
	Linear interpolation : Logarithmic interpolation :	ε(E) = 4.57 10 ⁻³ ε (E) = 1.87 10 ⁻³

Actual value (at 662 keV) : $\mathcal{E}(E) = 1.90 \ 10^{-3}$

Conclusion : approximation to be used only if high uncertainties (> 10 %) are acceptable

Efficiency calibration : mathematical fitting (1)

Determination of the best fitted function to a given set of experimental data (energy, efficiency)



In the logarithmic scale, the shape is smoother than in the linear scale.

Efficiency calibration : mathematical fitting

Functions frequently used:

Polynomial fitting in the log-log scale:

$$ln\varepsilon(E) = \sum_{i=0}^{n} a_i \cdot (lnE)^i$$

$$ln\varepsilon(E) = \sum_{i=0}^{n} a_i \cdot E^{-i}$$

Remarks :

-a, coefficients are determined using a least-squares fitting method

- experimental data must be weighted

- the polynomial degree (n) must be adjusted depending on the number of experimental data (p) : n << p

- in some case two different functions can be used with a cross point
- check the resulting fitted curves !

Efficiency calibration : mathematical fitting

Example : 40 experimental values in the 122-to-1836 keV range

Fitting function :

$$\ln \varepsilon(E) = \sum_{i=0}^{n} a_i \cdot (\ln E)^i$$

Adjusted coefficients :

fitting	122 to 1836 keV
deg0	-34,11961
deg1	48,16797
deg2	-25,89215
deg3	5,80219
deg4	-0,40503
deg5	-0,01632



Efficiency calibration : mathematical fitting

Efficiency fitting must be visually checked





In the case of cross points be carefull : - Avoid zones with important inflexion - Avoid high degree polynomials

Spline functions

Simply a curve

Special function defined piecewise by polynomials



- A piecewise polynomial f(x) is obtained by dividing of X into contiguous intervals, and representing f(x) by a separate polynomial in each interval
- The polynomials are joined together at the interval endpoints (knots) in such a way that a certain degree of smoothness of the resulting function is guaranteed

Piecewise defined polynomials



{K_i } (i = 1, ..., k+1)
K₁
$$\leq x_{min} < K_2 < ... < K_k < x_{max} < K_{k+1}$$

Series of knots

 x_{min} : minimum x-value of the data points x_{max} : maximum x-value of the data points

- 2 consecutive knots establish an interval [K_i, K_{i+1}) where a polynomial $P_i(x)$ of order n (degree n-1) is defined.

- In total k intervals are fixed by the knots.
- Outside the interval i the polynomial $P_i(x)$ is not defined.

• At the joining points $K_2 \dots K_k$ they have to fit smoothly – continous up to the (n-2) derivative

Spline functions mathematically

k polynomials in k intervals:

$$F(x) = \sum_{i=1}^{k} P_i(x)$$

where

$$P(x) = \begin{cases} \sum_{j=1}^{n} a_{ij} \cdot (x - K_i)^{j-1} & \text{for } K_i \le x \le K_{i+1} \\ 0 & \text{else} \end{cases}$$

Condition that a spline function solution must satisfy:

For the derivatons of neighboured polynomials at all knots:

$$P_i^{(m)}(K_{i+1}) - P_{i+1}^{(m)}(K_{i+1}) = 0, \quad 0 \le m \le n - 2$$

Efficiency calibration at 25 cm



Efficiency calibration at 10 cm



Multigamma emitters -> Coincidence summing effets

FEP Efficiency calibration : remarks

Efficiency calibration for reference geometry

 For point source : relative uncertainty 1-2 %

 Corrective factors needed if different measurement geometry

TOTAL EFFICIENCY

Total Efficiency (TE): $\mathcal{E}_T(E)$

Ratio of the total number of counts in the spectrum $(N_T(E))$, by the number of photons with energy *E* emitted by the source (F(E))



Calculation of the total efficiency

$$\mathcal{E}_{\mathcal{T}}(E) = \int_{\Omega} \prod_{i} \left(\exp\left(-\mu_{i}(E) \cdot x_{i}\right) \right) \cdot \left(1 - \exp\left(-\mu_{d}(E) \cdot x_{d}\right) \cdot d\Omega \right)$$

Attenuation in screens

Detector interaction

where x_i : i shield thickness x_d : detector thickness

Total efficiency depends on:

- Solid angle $\boldsymbol{\Omega}$
- Attenuations in absorbing layers (air, window, dead layer, etc)
- Interaction in the detector active volume (any effect)

And on the environment ! (scattering)

Experimental total efficiency calibration $\varepsilon_{T}(E)$



$$\mathcal{E}_{T}(E) = \frac{NT(E)}{A.I_{\gamma}(E)}$$

Experimental total efficiency calibration $\mathcal{E}_{T}(E)$

1. Using <u>mono-energetic</u> radionuclides

²⁴¹Am (60 keV), ¹⁰⁹Cd (88 keV), ¹³⁹Ce (166 keV- $T_{1/2}$ =138 d), ⁵¹Cr (320 keV- $T_{1/2}$ = 28d), ⁸⁵Sr (514 keV- $T_{1/2}$ =65 j), ¹³⁷Cs (662 keV), ⁵⁴Mn (834 keV - $T_{1/2}$ =302 d)

- 2. For the low-energy range, if P-type detector, ⁵⁷Co can be used (mean energy and sum of emission intensities)
- 3. For the high energy range, ⁸⁸Y and ⁶⁰Co can be used ... But complex decay scheme

Approximation of the total efficiency with 2-photons emitters: $N_T = A (I_{\gamma_1} \cdot \varepsilon_{T_1} + I_{\gamma_2} \cdot \varepsilon_{T_2} - I_{\gamma_1} \cdot \varepsilon_{T_1} \cdot P_{12} \cdot I_{\gamma_2} \cdot \varepsilon_{T_2})$ ⁸⁸Y: $E_1 = 898 \text{ keV}$ and $E_2 = 1836 \text{ keV}$ (or ⁶⁵Zn : 511 and 1115keV) $E_1 = 898 \text{ keV} \rightarrow \varepsilon_{T_1}$ can be extrapolated from previous data (⁵⁴Mn - 834 keV)

 $P_{12} = 1$



$$\varepsilon_{T2} = \frac{N_{\tau} - A \cdot I \gamma_{1} \cdot \varepsilon_{T1}}{A \cdot I_{\gamma 2} \cdot (1 - I_{\gamma 1} \cdot \varepsilon_{T1})}$$

Approximation of the total efficiency with 2-photons emitters:

$$N_{T} = A \left(I_{\gamma 1} \cdot \mathcal{E}_{T1} + I_{\gamma 2} \cdot \mathcal{E}_{T2} - I_{\gamma 1} \cdot \mathcal{E}_{T1} \cdot P_{12} \cdot I_{\gamma 2} \cdot \mathcal{E}_{T2} \right)$$



$$\varepsilon_{\tau} = \frac{N_{\tau} - A \cdot \varepsilon_{\tau}}{A \cdot (1 - \varepsilon_{\tau})} = 1 - \sqrt{1 - \frac{N_{\tau}}{A}}$$

Total Efficiency calibration : remarks

- Efficiency calibration for reference geometry

 Difficult, even with point sources:
 - relative uncertainty 5-10 %
 - But limited influence on the coincidence correction factor

• Monte Carlo simulation ?

MONTE CARLO SIMULATION

Monte Carlo simulation

- Difficulties: bad knowledge of the detector internal parameters
- Accurate description is time-consuming:
 - Radiography (external dimensions)
 - Collimated beam (hole dead layer)
 - Window to crystal distance (source at different distance)
 - Comparison with some experimental values
 - Different energies
 - Different geometries



X-ray tube -> Geometrical dimensions Cristal shape (rounding)



⁶⁰Co source -> Hole dimensions

Comparison with experimental data

Extended-range coaxial HPGe with carbon-epoxy window (61 x 61 mm) Point sources at 16 cm from the detector window



Picture from V. Peyres - CIEMAT

Monte Carlo simulation

• Can provide

- FEP efficiency
- Total efficiency (very dependent on the environment)
- Absolute calculation should be compared with accurate experimental data

Strong interest for efficiency transfer calculation